◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 三臣 - のへで

Open Questions in Open Systems

Claude-Alain Pillet

Centre de Physique Théorique (UMR 6207) Université de Toulon

ANR HAM-MARK 2009

1 Hamiltonian and markovian models

2 Where we stand (a rough picture)

3 Where to go



◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

Open quantum systems



(ロ)、(型)、(E)、(E)、 E) のQの

Hamiltonian models

• C^* - or W^* -algebra \mathfrak{A} generated by sub-algebras \mathfrak{A}_S , $\mathfrak{A}_{\mathcal{R}_1}$, ...

- C^* or W^* -algebra \mathfrak{A} generated by sub-algebras \mathfrak{A}_S , $\mathfrak{A}_{\mathcal{R}_1}$, ...
- Groups $\tau_{\mathcal{S}}^t = e^{t\delta_{\mathcal{S}}}, \ \tau_{\mathcal{R}_1}^t = e^{t\delta_{\mathcal{R}_1}}, \ldots$ of *-automorphisms of $\mathfrak{A}_{\mathcal{S}}, \ \mathfrak{A}_{\mathcal{R}_1}, \ldots$

- C^* or W^* -algebra \mathfrak{A} generated by sub-algebras \mathfrak{A}_S , $\mathfrak{A}_{\mathcal{R}_1}$, ...
- Groups $\tau_{\mathcal{S}}^t = e^{t\delta_{\mathcal{S}}}$, $\tau_{\mathcal{R}_1}^t = e^{t\delta_{\mathcal{R}_1}}$, ... of *-automorphisms of $\mathfrak{A}_{\mathcal{S}}$, $\mathfrak{A}_{\mathcal{R}_1}$, ...
- Reference state ω of \mathfrak{A} , e.g., $\omega|_{\mathcal{R}_j}$ is $\tau^t_{\mathcal{R}_i}$ -KMS state.

- C^* or W^* -algebra \mathfrak{A} generated by sub-algebras \mathfrak{A}_S , $\mathfrak{A}_{\mathcal{R}_1}$, ...
- Groups $\tau_{\mathcal{S}}^t = e^{t\delta_{\mathcal{S}}}$, $\tau_{\mathcal{R}_1}^t = e^{t\delta_{\mathcal{R}_1}}$, ... of *-automorphisms of $\mathfrak{A}_{\mathcal{S}}$, $\mathfrak{A}_{\mathcal{R}_1}$, ...
- Reference state ω of \mathfrak{A} , e.g., $\omega|_{\mathcal{R}_i}$ is $\tau^t_{\mathcal{R}_i}$ -KMS state.
- Coupling $V_j = V_j^* \in \mathfrak{A}_S \lor \mathfrak{A}_{\mathcal{R}_j}$.

- C^* or W^* -algebra \mathfrak{A} generated by sub-algebras \mathfrak{A}_S , $\mathfrak{A}_{\mathcal{R}_1}$, ...
- Groups $\tau_{\mathcal{S}}^t = e^{t\delta_{\mathcal{S}}}$, $\tau_{\mathcal{R}_1}^t = e^{t\delta_{\mathcal{R}_1}}$, ... of *-automorphisms of $\mathfrak{A}_{\mathcal{S}}$, $\mathfrak{A}_{\mathcal{R}_1}$, ...
- Reference state ω of \mathfrak{A} , e.g., $\omega|_{\mathcal{R}_i}$ is $\tau^t_{\mathcal{R}_i}$ -KMS state.
- Coupling $V_j = V_j^* \in \mathfrak{A}_S \vee \mathfrak{A}_{\mathcal{R}_j}$.
- Coupled dynamics $\tau^t = e^{t\delta}$, $\delta = \delta_S + \delta_{R_1} + \cdots + i[V_1, \cdot] + \cdots$

- C^* or W^* -algebra \mathfrak{A} generated by sub-algebras \mathfrak{A}_S , $\mathfrak{A}_{\mathcal{R}_1}$, ...
- Groups $\tau_{\mathcal{S}}^t = e^{t\delta_{\mathcal{S}}}, \ \tau_{\mathcal{R}_1}^t = e^{t\delta_{\mathcal{R}_1}}, \ldots$ of *-automorphisms of $\mathfrak{A}_{\mathcal{S}}, \ \mathfrak{A}_{\mathcal{R}_1}, \ldots$
- Reference state ω of \mathfrak{A} , e.g., $\omega|_{\mathcal{R}_j}$ is $\tau^t_{\mathcal{R}_i}$ -KMS state.
- Coupling $V_j = V_j^* \in \mathfrak{A}_S \lor \mathfrak{A}_{\mathcal{R}_j}$.
- Coupled dynamics $\tau^t = e^{t\delta}$, $\delta = \delta_S + \delta_{R_1} + \cdots + i[V_1, \cdot] + \cdots$
- Natural Steady State

$$\omega^+ = \mathbf{w}^*_{t \to \infty} - \lim_{t \to \infty} \frac{1}{t} \int_0^t \omega \circ \tau^s \, \mathrm{d}s$$

Hamiltonian models

- C^* or W^* -algebra \mathfrak{A} generated by sub-algebras \mathfrak{A}_S , $\mathfrak{A}_{\mathcal{R}_1}$, ...
- Groups $\tau_{\mathcal{S}}^t = e^{t\delta_{\mathcal{S}}}$, $\tau_{\mathcal{R}_1}^t = e^{t\delta_{\mathcal{R}_1}}$, ... of *-automorphisms of $\mathfrak{A}_{\mathcal{S}}$, $\mathfrak{A}_{\mathcal{R}_1}$, ...
- Reference state ω of \mathfrak{A} , e.g., $\omega|_{\mathcal{R}_j}$ is $\tau^t_{\mathcal{R}_i}$ -KMS state.
- Coupling $V_j = V_j^* \in \mathfrak{A}_S \vee \mathfrak{A}_{\mathcal{R}_j}$.
- Coupled dynamics $\tau^t = e^{t\delta}$, $\delta = \delta_S + \delta_{R_1} + \cdots + i[V_1, \cdot] + \cdots$
- Natural Steady State

$$\omega^+ = w^*_{t\to\infty} - \lim_{t\to\infty} \frac{1}{t} \int_0^t \omega \circ \tau^s \, \mathrm{d}s.$$

Mean entropy production rate

$$\operatorname{Ep}(\omega^+) = -\lim_{t\to\infty} \frac{1}{t}\operatorname{Ent}(\omega\circ\tau^t|\omega).$$

Markovian models: Level I

• "Small" W^* -algebra $\mathfrak{A}_S = \mathcal{B}(\mathcal{H}_S)$.



◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

Markovian models: Level I

- "Small" W^* -algebra $\mathfrak{A}_S = \mathcal{B}(\mathcal{H}_S)$.
- $\mathfrak{A}_{\mathcal{S}*}$ identified with trace class operators on $\mathcal{H}_{\mathcal{S}}$.

Markovian models: Level I

- "Small" W^* -algebra $\mathfrak{A}_S = \mathcal{B}(\mathcal{H}_S)$.
- $\mathfrak{A}_{\mathcal{S}*}$ identified with trace class operators on $\mathcal{H}_{\mathcal{S}}$.
- Quantum master equation in Heisenberg picture

$$\frac{\mathrm{d}}{\mathrm{d}t}A = \mathcal{L}(A) = \mathrm{i}[H, A] + \frac{1}{2}\sum_{\mathsf{L}}V_{\mathsf{L}}^*[A, V_{\mathsf{L}}] + [V_{\mathsf{L}}^*, A]V_{\mathsf{L}}.$$

Markovian models: Level I

- "Small" W^* -algebra $\mathfrak{A}_S = \mathcal{B}(\mathcal{H}_S)$.
- $\mathfrak{A}_{\mathcal{S}*}$ identified with trace class operators on $\mathcal{H}_{\mathcal{S}}$.
- Quantum master equation in Heisenberg picture

$$\frac{\mathrm{d}}{\mathrm{d}t}A = \mathcal{L}(A) = \mathrm{i}[H,A] + \frac{1}{2}\sum_{\iota}V_{\iota}^*[A,V_{\iota}] + [V_{\iota}^*,A]V_{\iota}.$$

• Semi-group $\alpha^t = e^{t\mathcal{L}}$ of CP, 1-preserving maps on $\mathfrak{A}_{\mathcal{S}}$.

Markovian models: Level I

- "Small" W^* -algebra $\mathfrak{A}_S = \mathcal{B}(\mathcal{H}_S)$.
- $\mathfrak{A}_{\mathcal{S}*}$ identified with trace class operators on $\mathcal{H}_{\mathcal{S}}$.
- Quantum master equation in Heisenberg picture

$$\frac{\mathrm{d}}{\mathrm{d}t}A = \mathcal{L}(A) = \mathrm{i}[H, A] + \frac{1}{2}\sum_{\iota} V_{\iota}^*[A, V_{\iota}] + [V_{\iota}^*, A]V_{\iota}.$$

- Semi-group $\alpha^t = e^{t\mathcal{L}}$ of CP, 1-preserving maps on $\mathfrak{A}_{\mathcal{S}}$.
- Steady state condition $\mathcal{L}^*(\rho) = 0$.

Markovian models: Level I

- "Small" W^* -algebra $\mathfrak{A}_S = \mathcal{B}(\mathcal{H}_S)$.
- $\mathfrak{A}_{\mathcal{S}*}$ identified with trace class operators on $\mathcal{H}_{\mathcal{S}}$.
- Quantum master equation in Heisenberg picture

$$\frac{\mathrm{d}}{\mathrm{d}t}A = \mathcal{L}(A) = \mathrm{i}[H, A] + \frac{1}{2}\sum_{\iota} V_{\iota}^*[A, V_{\iota}] + [V_{\iota}^*, A]V_{\iota}.$$

- Semi-group $\alpha^t = e^{t\mathcal{L}}$ of CP, 1-preserving maps on $\mathfrak{A}_{\mathcal{S}}$.
- Steady state condition $\mathcal{L}^*(\rho) = 0$.
- Convergence to steady state

$$w - \lim_{t \to \infty} e^{t\mathcal{L}^*}(\omega) = \rho$$

Markovian models: Level II

•
$$\mathcal{Z} = \mathcal{H}_{\mathcal{S}} \otimes \Gamma_{+}(L^{2}(\mathbb{R}, d\tau) \otimes \mathfrak{h})$$



Markovian models: Level II

- $\mathcal{Z} = \mathcal{H}_{\mathcal{S}} \otimes \Gamma_{+}(L^{2}(\mathbb{R}, d\tau) \otimes \mathfrak{h})$
- Quantum Langevin equation in Schrödinger picture

$$\mathrm{i}\frac{\mathrm{d}}{\mathrm{d}t}U^{t} = (H \otimes I - I \otimes \mathrm{d}\Gamma(\mathrm{i}\partial_{\tau}))U^{t} + \sum_{\iota} V_{\iota} \otimes a_{\iota}^{*}(\delta_{0})U^{t} + V_{\iota}^{*} \otimes a_{\iota}(\delta_{0})U^{t},$$

generates a unitary group U^t on \mathcal{Z} .

Markovian models: Level II

- $\mathcal{Z} = \mathcal{H}_{\mathcal{S}} \otimes \Gamma_{+}(L^{2}(\mathbb{R}, d\tau) \otimes \mathfrak{h})$
- Quantum Langevin equation in Schrödinger picture

$$\mathrm{i}\frac{\mathrm{d}}{\mathrm{d}t}U^{t} = (H \otimes I - I \otimes \mathrm{d}\Gamma(\mathrm{i}\partial_{\tau}))U^{t} + \sum_{\iota} V_{\iota} \otimes a_{\iota}^{*}(\delta_{0})U^{t} + V_{\iota}^{*} \otimes a_{\iota}(\delta_{0})U^{t},$$

generates a unitary group U^t on \mathcal{Z} .

Associated Markovian semigroup

$$e^{t\mathcal{L}}(A) = \langle \Omega | U^t(A \otimes I) U^{t*} | \Omega \rangle$$

is generated by

$$\mathcal{L}(A) = \mathrm{i}[H, A] + \frac{1}{2} \sum_{\iota} V_{\iota}^*[A, V_{\iota}] + [V_{\iota}^*, A] V_{\iota}.$$

・ロト ・ 雪 ト ・ ヨ ト

э

Where to go

Repeated interaction models



◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

Markovian models

• Return to equilibrium [Hepp-Lieb, Davies in the early 70'].

- Return to equilibrium [Hepp-Lieb, Davies in the early 70'].
- Convergence to steady state and nonequilibrium thermodynamics [Davies, Davies-Spohn, Lebowitz-Spohn in the late 70'].

- Return to equilibrium [Hepp-Lieb, Davies in the early 70'].
- Convergence to steady state and nonequilibrium thermodynamics [Davies, Davies-Spohn, Lebowitz-Spohn in the late 70'].
- Repeated interactions \longrightarrow quantum Langevin [Attal-Pautrat, Attal-Joye].

- Return to equilibrium [Hepp-Lieb, Davies in the early 70'].
- Convergence to steady state and nonequilibrium thermodynamics [Davies, Davies-Spohn, Lebowitz-Spohn in the late 70'].
- Repeated interactions \longrightarrow quantum Langevin [Attal-Pautrat, Attal-Joye].
- Repeated interaction [Bruneau-Joye-Merkli, Bruneau-P].

- Return to equilibrium [Hepp-Lieb, Davies in the early 70'].
- Convergence to steady state and nonequilibrium thermodynamics [Davies, Davies-Spohn, Lebowitz-Spohn in the late 70'].
- Repeated interactions \longrightarrow quantum Langevin [Attal-Pautrat, Attal-Joye].
- Repeated interaction [Bruneau-Joye-Merkli, Bruneau-P].
- Current fluctuations [Avron-Bachmann-Graf-Klich, De Roeck-Dereziński-Maes].

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

$\mathsf{Hamiltonian}\longleftrightarrow\mathsf{Markovian}$

• Weak coupling (van Hove) limit [Davies 70'].

Outline

$\mathsf{Hamiltonian}\longleftrightarrow\mathsf{Markovian}$

- Weak coupling (van Hove) limit [Davies 70'].
- Unitary dilation and quantum stochastic evolution [Accardi-Frigerio-Lewis, Hudson-Parthasarathy, Frigerio, Maassen mid 80'].

Outline

$\mathsf{Hamiltonian}\longleftrightarrow\mathsf{Markovian}$

- Weak coupling (van Hove) limit [Davies 70'].
- Unitary dilation and quantum stochastic evolution [Accardi-Frigerio-Lewis, Hudson-Parthasarathy, Frigerio, Maassen mid 80'].
- Stochastic limit [Accardi-Frigerio-Lu-Volovich 90', 00'].

$\mathsf{Hamiltonian} \longleftrightarrow \mathsf{Markovian}$

- Weak coupling (van Hove) limit [Davies 70'].
- Unitary dilation and quantum stochastic evolution [Accardi-Frigerio-Lewis, Hudson-Parthasarathy, Frigerio, Maassen mid 80'].
- Stochastic limit [Accardi-Frigerio-Lu-Volovich 90', 00'].
- Extended weak coupling limit [De Roeck-Dereziński].

$\mathsf{Hamiltonian}\longleftrightarrow\mathsf{Markovian}$

- Weak coupling (van Hove) limit [Davies 70'].
- Unitary dilation and quantum stochastic evolution [Accardi-Frigerio-Lewis, Hudson-Parthasarathy, Frigerio, Maassen mid 80'].
- Stochastic limit [Accardi-Frigerio-Lu-Volovich 90', 00'].
- Extended weak coupling limit [De Roeck-Dereziński].
- Generalized weak coupling limit [Taj].

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

Hamiltonian models

 Return to equilibrium [Robinson 70', Jakšić-P, Dereziński-Jakšić, Bach-Fröhlich-Sigal, Fröhlich-Merkli].

- Return to equilibrium [Robinson 70', Jakšić-P, Dereziński-Jakšić, Bach-Fröhlich-Sigal, Fröhlich-Merkli].
- Convergence to steady state and nonequilibrium thermodynamics [Jakšić-P, Fröhlich-Merkli-Ueltschi, Abou Salem-Fröhlich, Merkli-Mück-Sigal].

- Return to equilibrium [Robinson 70', Jakšić-P, Dereziński-Jakšić, Bach-Fröhlich-Sigal, Fröhlich-Merkli].
- Convergence to steady state and nonequilibrium thermodynamics [Jakšić-P, Fröhlich-Merkli-Ueltschi, Abou Salem-Fröhlich, Merkli-Mück-Sigal].
- Fluctuation-dissipation near equilibrium [Jakšić-Ogata-Pautrat-P].



 MMS combine the L[∞]-Liouvillian approach with a hybrid spectral deformation technique to prove existence of NESS in open systems coupled to bosonic reservoirs



- MMS combine the L[∞]-Liouvillian approach with a hybrid spectral deformation technique to prove existence of NESS in open systems coupled to bosonic reservoirs
- The MMS machinery is quite heavy.



- MMS combine the L[∞]-Liouvillian approach with a hybrid spectral deformation technique to prove existence of NESS in open systems coupled to bosonic reservoirs
- The MMS machinery is quite heavy.
- A streamlined approach to the problem could have very important consequences in the field (somewhat like the Aizenman-Molchanov proof of Anderson localization).



- MMS combine the L[∞]-Liouvillian approach with a hybrid spectral deformation technique to prove existence of NESS in open systems coupled to bosonic reservoirs
- The MMS machinery is quite heavy.
- A streamlined approach to the problem could have very important consequences in the field (somewhat like the Aizenman-Molchanov proof of Anderson localization).
- A starting point may be a more clever choice of the cyclic vector in the standard representation of \mathfrak{A} .

The strong coupling limit

• Except for quasi-free systems where scattering theory plays the central rôle, all available techniques in open systems are perturbative:

- Except for quasi-free systems where scattering theory plays the central rôle, all available techniques in open systems are perturbative:
 - in the coupling to the environment;

- Except for quasi-free systems where scattering theory plays the central rôle, all available techniques in open systems are perturbative:
 - in the coupling to the environment;
 - in the (non-quadratic) interactions.

- Except for quasi-free systems where scattering theory plays the central rôle, all available techniques in open systems are perturbative:
 - in the coupling to the environment;
 - in the (non-quadratic) interactions.
- Many important physical problems are beyond reach of these perturbative approaches.

- Except for quasi-free systems where scattering theory plays the central rôle, all available techniques in open systems are perturbative:
 - in the coupling to the environment;
 - in the (non-quadratic) interactions.
- Many important physical problems are beyond reach of these perturbative approaches.
- Exact solution, via Bethe Ansatz ?

・ロト ・ 一 ト ・ モト ・ モト

э.

Where to go

The strong coupling limit



・ロト ・四ト ・ヨト ・ヨト

э

Where to go

The strong coupling limit



$$H_L = -t \sum_{x < 0,\sigma} a^*_{\sigma}(x) a_{\sigma}(x-1) + a^*_{\sigma}(x-1) a_{\sigma}(x)$$

ヘロン 人間 とくほとう ほとう

э

Where to go

The strong coupling limit



$$H_{R} = -t \sum_{x > 0,\sigma} a_{\sigma}^{*}(x) a_{\sigma}(x+1) + a_{\sigma}^{*}(x+1) a_{\sigma}(x)$$

ヘロン 人間 とくほとう ほとう

э

Where to go

The strong coupling limit



$$H_{D} = -\sum_{\sigma} \varepsilon_{0} a_{\sigma}^{*}(0) a_{\sigma}(0) + U a_{-}^{*}(0) a_{-}(0) a_{+}^{*}(0) a_{+}(0)$$

・ロト ・ 一 ト ・ モト ・ モト

э.

Where to go





$$H = H_L + H_R + H_D + \lambda \sum_{x \in \{\pm 1\}, \sigma} a^*_{\sigma}(x) a_{\sigma}(0) + a^*_{\sigma}(0) a_{\sigma}(x)$$

・ロト ・ 理 ト ・ ヨ ト ・ ヨ ト

э

The strong coupling limit



The Kondo regime

 $U \gg \lambda \simeq t$

is non-perturbative.

イロト 不得 トイヨト イヨト

э.

Where to go

The strong coupling limit



Is there a "non-linear" Landauer-Büttiker formula ?

Large systems: Hamiltonian approach

Except again for quasi-free systems, all available techniques in open systems are limited to "small" system S with finitely many degrees of freedom (dim H_S < ∞), e.g.,

- Except again for quasi-free systems, all available techniques in open systems are limited to "small" system S with finitely many degrees of freedom (dim H_S < ∞), e.g.,
 - in matter-radiation systems, atoms have a finite number of levels;

- Except again for quasi-free systems, all available techniques in open systems are limited to "small" system S with finitely many degrees of freedom (dim H_S < ∞), e.g.,
 - in matter-radiation systems, atoms have a finite number of levels;
 - in mesoscopic systems multibody interactions can only take place on a finite dimensional subspace of the 1-body Hilbert space.

- Except again for quasi-free systems, all available techniques in open systems are limited to "small" system S with finitely many degrees of freedom (dim H_S < ∞), e.g.,
 - in matter-radiation systems, atoms have a finite number of levels;
 - in mesoscopic systems multibody interactions can only take place on a finite dimensional subspace of the 1-body Hilbert space.
- This limitation is mainly due to the use of perturbative techniques:

- Except again for quasi-free systems, all available techniques in open systems are limited to "small" system S with finitely many degrees of freedom (dim H_S < ∞), e.g.,
 - in matter-radiation systems, atoms have a finite number of levels;
 - in mesoscopic systems multibody interactions can only take place on a finite dimensional subspace of the 1-body Hilbert space.
- This limitation is mainly due to the use of perturbative techniques:
 - in the Liouvillian approach, the Fermi golden rule only gives uniform control of finitely many resonances;

- Except again for quasi-free systems, all available techniques in open systems are limited to "small" system S with finitely many degrees of freedom (dim H_S < ∞), e.g.,
 - in matter-radiation systems, atoms have a finite number of levels;
 - in mesoscopic systems multibody interactions can only take place on a finite dimensional subspace of the 1-body Hilbert space.
- This limitation is mainly due to the use of perturbative techniques:
 - in the Liouvillian approach, the Fermi golden rule only gives uniform control of finitely many resonances;
 - in the C*-scattering approach, the Dyson expansion can not be controlled uniformly in the "dimension" of the perturbation.

Large systems: Markovian approach

• The situation is better in the Markovian approach. Some results for repeated interaction systems with infinite dimensional "small" system:

- The situation is better in the Markovian approach. Some results for repeated interaction systems with infinite dimensional "small" system:
 - repeated interactions of a 1-mode QED cavity with 2-level atoms;

- The situation is better in the Markovian approach. Some results for repeated interaction systems with infinite dimensional "small" system:
 - repeated interactions of a 1-mode QED cavity with 2-level atoms;
 - repeated interactions of a particle on a tight-binding lattice with 2-level atoms.

- The situation is better in the Markovian approach. Some results for repeated interaction systems with infinite dimensional "small" system:
 - repeated interactions of a 1-mode QED cavity with 2-level atoms;
 - repeated interactions of a particle on a tight-binding lattice with 2-level atoms.
- The usual Davies approach to the weak coupling limit also suffer from the restriction to finite dimensional "small" system. A fact which strongly limit the use of Markovian description of such systems. Some recent results by [Taj] may change this! (continuous spectrum)

- The situation is better in the Markovian approach. Some results for repeated interaction systems with infinite dimensional "small" system:
 - repeated interactions of a 1-mode QED cavity with 2-level atoms;
 - repeated interactions of a particle on a tight-binding lattice with 2-level atoms.
- The usual Davies approach to the weak coupling limit also suffer from the restriction to finite dimensional "small" system. A fact which strongly limit the use of Markovian description of such systems. Some recent results by [Taj] may change this! (continuous spectrum)
- Quantum Brownian motion.

(ロ)、(型)、(E)、(E)、 E) の(の)

The rotating wave approximation

• Matter-radiation problem

$$H = \omega_0 b^* b + \mathrm{d}\Gamma(|k|) + \lambda \left(b + b^*\right) \left(a(f) + a^*(f)\right).$$

The rotating wave approximation

• Matter-radiation problem

$$H = \omega_0 b^* b + \mathrm{d}\Gamma(|k|) + \lambda \left(b + b^*\right) \left(a(f) + a^*(f)\right).$$

• Resonant interaction: $||k| - \omega_0| \ll \omega_0$ on supp f.

・ロト・日本・モート モー うへで

The rotating wave approximation

• Matter-radiation problem

$$H = \omega_0 b^* b + \mathrm{d}\Gamma(|k|) + \lambda \left(b + b^*\right) \left(a(f) + a^*(f)\right).$$

- Resonant interaction: $||k| \omega_0| \ll \omega_0$ on supp f.
- RW approximation

$$H = \omega_0 b^* b + \mathrm{d}\Gamma(|k|) + \lambda \left(b^* a(f) + b a^*(f) \right).$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

The rotating wave approximation

Matter-radiation problem

$$H = \omega_0 b^* b + \mathrm{d}\Gamma(|k|) + \lambda \left(b + b^*\right) \left(a(f) + a^*(f)\right).$$

- Resonant interaction: $||k| \omega_0| \ll \omega_0$ on supp f.
- RW approximation

$$H = \omega_0 b^* b + \mathrm{d}\Gamma(|k|) + \lambda \left(b^* a(f) + b a^*(f) \right).$$

• Get effective bounds on dynamics at positive temperature.

The rotating wave approximation

Matter-radiation problem

$$H = \omega_0 b^* b + \mathrm{d}\Gamma(|k|) + \lambda \left(b + b^*\right) \left(a(f) + a^*(f)\right).$$

- Resonant interaction: $||k| \omega_0| \ll \omega_0$ on supp f.
- RW approximation

$$H = \omega_0 b^* b + \mathrm{d}\Gamma(|k|) + \lambda \left(b^* a(f) + b a^*(f) \right).$$

- Get effective bounds on dynamics at positive temperature.
- Compare the corresponding Markovian approximations.

The rotating wave approximation

Matter-radiation problem

$$H = \omega_0 b^* b + \mathrm{d}\Gamma(|k|) + \lambda \left(b + b^*\right) \left(a(f) + a^*(f)\right).$$

- Resonant interaction: $||k| \omega_0| \ll \omega_0$ on supp f.
- RW approximation

$$H = \omega_0 b^* b + \mathrm{d}\Gamma(|k|) + \lambda \left(b^* a(f) + b a^*(f) \right).$$

- Get effective bounds on dynamics at positive temperature.
- Compare the corresponding Markovian approximations.
- Apply the RW approximation to matter-radiation problems and do the relevant spectral analysis (generalization of Friedrichs type models).

Quantum probability & statistics

• Quantum CLT for sums of iid variables (extends Kuperberg's result to non-tracial states).

Quantum probability & statistics

- Quantum CLT for sums of iid variables (extends Kuperberg's result to non-tracial states).
- Quantum hypothesis testing and nonequilibrium thermodynamics.

Quantum probability & statistics

- Quantum CLT for sums of iid variables (extends Kuperberg's result to non-tracial states).
- Quantum hypothesis testing and nonequilibrium thermodynamics.
- Quantum random walks.

Quantum probability & statistics

- Quantum CLT for sums of iid variables (extends Kuperberg's result to non-tracial states).
- Quantum hypothesis testing and nonequilibrium thermodynamics.
- Quantum random walks.
- . . .