

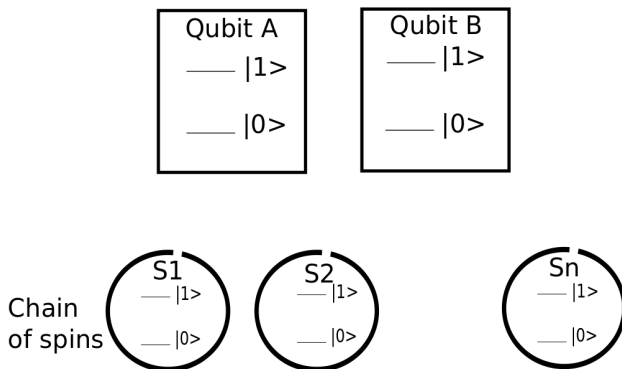
Evolution of entanglement of two qubits with a model of repeated interactions

Sylvain Vogelsberger

Université Joseph Fourier, Institut Fourier, Grenoble (France)

29th September 2010 Grenoble

Model : 2 qubits in contact with a heat bath consisting in an infinite chain of spins at thermal equilibrium at inverse temperature β .



More precisely :

- $h_0^{qubits} = \epsilon[\sigma_+ \sigma_-^{(A)} + \sigma_+ \sigma_-^{(B)}]$ is the free hamiltonian for qubits.
- $h_0^{chain} = \sum_{k=1} \epsilon a^\dagger a(k)$ is the free hamiltonian for the chain.
- During time $[(k-1)\tau, (k-1)\tau + \frac{\tau}{2}]$, qubit A interacts with spin S_k for $k \in \mathbb{N}$. The interaction is given by
$$W^{(A)} = \lambda \left[\sigma_+^{(A)} \otimes a(k) + \sigma_-^{(A)} \otimes a^\dagger(k) \right].$$
- Then, qubit B interacts with the same spin S_k during time $[(k-1)\tau + \frac{\tau}{2}, k\tau]$ with a similar interaction $W^{(B)}$.

Weak coupling limit : As Attal and Joye in Journal of Statistical Physics , **126** (2007), we take the limit :

- $T = \tau \frac{t}{\lambda^2}$ with $\frac{t}{\lambda^2} \in \mathbb{N}$ is the number of interactions and T is the macroscopic time scale
- $\frac{\tau}{2} = 1$
- $\lambda \rightarrow 0$

Initial state : We consider that at time $T = 0$, the system of two qubits is initially in an X-state $\rho(0)$, represented in the basis $\{|11\rangle, |10\rangle, |01\rangle, |00\rangle\}$ by :

$$\rho(0) = \begin{pmatrix} a & 0 & 0 & y \\ 0 & b & x & 0 \\ 0 & \bar{x} & c & 0 \\ \bar{y} & 0 & 0 & d \end{pmatrix} \quad (1)$$

Furthermore, we suppose that the chain of spins is initially in a thermal state at inverse temperature β .

In the limit above and with $\beta = \infty$, we have asymptotically ($T = \tau \frac{t}{\lambda^2} \rightarrow \infty$) :

$$\rho(t) = \begin{pmatrix} a(t) & 0 & 0 & y(t) \\ 0 & b(t) & x(t) & 0 \\ 0 & x(\bar{t}) & c(t) & 0 \\ y(\bar{t}) & 0 & 0 & d(t) \end{pmatrix} \quad (2)$$

with :

- $a(t) = e^{-2t}a(0)$
- $b(t) = [-e^{-2t} + e^{-t}]a(0) + e^{-t}b(0)$
- $c(t) = [-5e^{-2t} + (5 - 4t + t^2)e^{-t}]a(0) + t^2e^{-t}b(0) + e^{-t}c(0) + (-te^{-t}(x + \bar{x})(0))$
- $d(t) = [5e^{-2t} + (-6 + 4t - t^2)e^{-t} + 1]a(0) + [(-1 - t^2)e^{-t} + 1]b(0) + (1 - e^{-t})c(0) + d(0) + te^{-t}(x + \bar{x})(0)$
- $|y(t)| = e^{-t}|y(0)|$
- $|x(t)| =$
 $|(-2e^{-2t} + (2 - t)e^{-t})a(0) + (-te^{-t})b(0) + e^{-t}x(0)|$

In the case of $\beta = 0$, we have the same structure for the density matrix but :

- $a(t) = s_1(t)a(0) + s_3(t)[b(0) + c(0) + (x + \bar{x})(0)] + s_2(t)d(0)$
- $b(t) =$
 $s_3(t)a(0) + s_1(t)b(0) - s_3(t)(x + \bar{x})(0)] + s_2(t)c(0) + s_3d(0)$
- $c(t) =$
 $s_3(t)a(0) + s_2(t)b(0) - s_3(t)(x + \bar{x})(0)] + s_1(t)c(0) + s_3d(0)$
- $d(t) = s_2(t)a(0) + s_3(t)[b(0) + c(0) + (x + \bar{x})(0)] + s_1d(0)$
- $|x(t)| =$
 $|s_3(t)[a(0) + d(0) - (b(0) + c(0))] + s_1(t)x(0) + s_2(t)\bar{x}(0)|$
- $|y(t)| = e^{-t}|y(0)|$

with :

- $s_1(t) = \frac{1}{3} + \frac{1}{2}e^{-t} + \frac{1}{6}e^{-3t}$
- $s_2(t) = \frac{1}{3} - \frac{1}{2}e^{-t} + \frac{1}{6}e^{-3t}$
- $s_3(t) = \frac{1}{6}(1 - e^{-3t})$

Concurrence for an X-state : We measure entanglement via the concurrence C of Wootters. In the case of an X-state, we have a beautiful formula for the concurrence :

$$C(\rho) = 2 \max\{0, |x| - \sqrt{ad}, |y| - \sqrt{bc}\} \quad (3)$$

So it is easy to calculate...

we have done it and we find the following results :

for $\beta = \infty$

We can vary the parameter t and see the evolution of $C(t)$:

Proposition

For $\beta = \infty$ and for an initial X-state

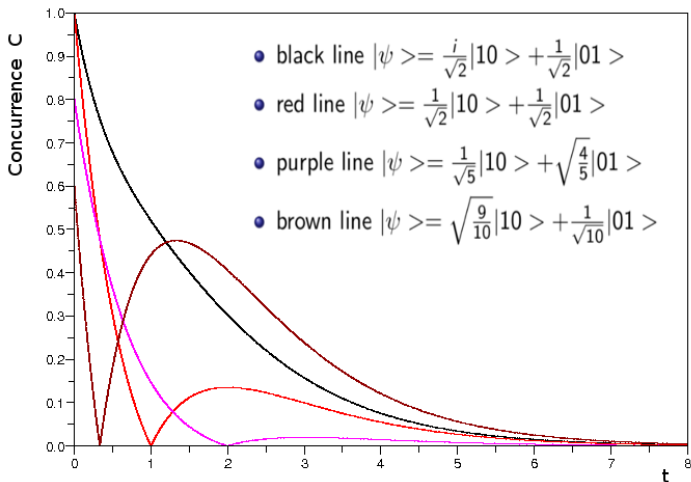
- *if $a + b > 0$ then $C(t) \underset{t \rightarrow \infty}{\sim} 2(a + b)te^{-t}$.*
- *Else $C(t) \underset{t \rightarrow \infty}{\sim} 2 \max\{0, |x|, |y|\}e^{-t}$ (here $(x, y) \neq (0, 0)$).*

A curious phenomenon :

Proposition

We suppose that initially $\rho(0) = |\psi\rangle\langle\psi|$ where $|\psi\rangle = \alpha|10\rangle + \beta|01\rangle$ then

- if $\arg \alpha = \arg \beta$ then $C(\frac{\beta}{\alpha}) = 0$ and $C(t) > 0$ for $t \neq \frac{\beta}{\alpha}$.*
- else $C(t) > 0$ for all t .*

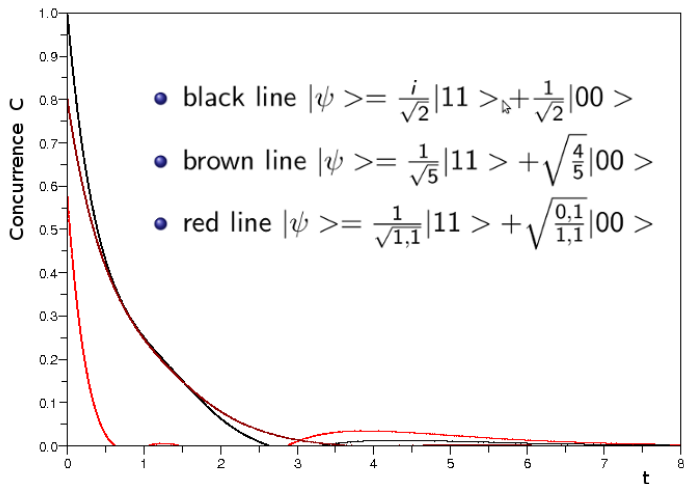


An other :

Proposition

We suppose that initially $\rho(0) = |\psi\rangle\langle\psi|$ where $|\psi\rangle = \alpha|11\rangle + \beta|00\rangle$ then there exists two constants $c_1 \simeq 1.093$ and $c_2 \simeq 1.106$ such as :

- if $c_1 \leq \frac{1}{|\alpha|^2} \leq c_2$ then there exists two steps such as $C(t) = 0$.*
- else there exists only one step such as $C(t) = 0$.*



For $\beta = 0$

Proposition

For $\beta = 0$ and for an initial X -state

- *if $(b + c) - (x + \bar{x}) < 1$, there exists $t_0 < \infty$ such as $C(t) = 0$ for all $t \geq t_0$.*
- *if $(b + c) - (x + \bar{x}) > 1$ then $C(t \rightarrow \infty) = (b + c) - (x + \bar{x}) - 1 > 0$.*
- *if $\begin{cases} (b+c)-(x+\bar{x})=1 \\ a=d \text{ and } \Im(x)=0 \end{cases}$ there exists $t_0 < \infty$ such as $C(t) = 0$ for all $t \geq t_0$.*
- *if $\begin{cases} (b+c)-(x+\bar{x})=1 \\ a \neq d \text{ or } \Im(x) \neq 0 \end{cases}$ then $C(t) \underset{t \rightarrow \infty}{\sim} 3[(\Im x)^2 + \frac{1}{2}(a - d)^2]e^{-2t}$.*

A bit more on $\mathcal{E}_{\mathcal{X}}$, the set of X-states :

Proposition

$\mathcal{E}_{\mathcal{X}}$ is a convex compact set of 7 real dimensions.

Due to Krein-Milman theorem, we know that $\mathcal{E}_{\mathcal{X}}$ is the convex hull of its extreme points.

Proposition

The extreme points of $\mathcal{E}_{\mathcal{X}}$ are the Bell states. Id est the states :

- $|\psi\rangle\langle\psi|$ for $|\psi\rangle = \alpha|11\rangle + \delta|00\rangle$ ($\alpha, \delta \in \mathbb{C}$)
- $|\phi\rangle\langle\phi|$ for $|\phi\rangle = \beta|10\rangle + \gamma|01\rangle$ ($\beta, \gamma \in \mathbb{C}$)

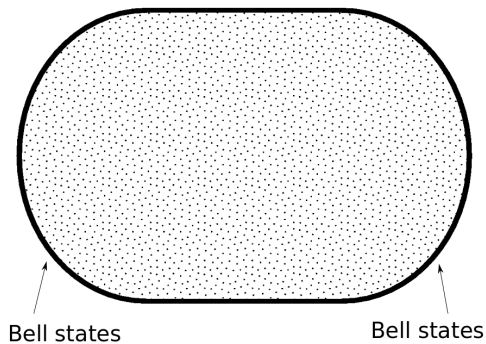



Figure: A view of \mathcal{E}_X

Proposition

The concurrence function, $C : \mathcal{E}_{\mathcal{X}} \rightarrow [0,1]$ is convex.

-  S. Attal and A. Joye, *Journal of Statistical Physics* , **126** (2007)
-  W.K. Wootters, *Phys. Rev. Lett.* **80**, 2245 (1998)
-  H.-P. Breuer and F. Petruccione, *The Theory of Open Quantum Systems* (Oxford University Press, 2002)
-  S. Haroche and J.-M. Raimond, *Exploring the quantum: atoms, cavities and photons* (Oxford Univ. Press, 2006)