On the scattering theory of asymptotically flat manifolds and Strichartz inequalities

Jean-Marc Bouclet Institut de Mathématiques de Toulouse

23 Juin 2016 - Cergy Pontoise

Conférence en l'honneur de Vladimir Georgescu

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

Introduction

Purpose of the talk

- Take the question of Strichartz inequalities (for the Schrödinger equation) on asymptotically flat manifolds as a case study to review some related scattering estimates (resolvent estimates, timedecay, smoothing estimates), either for comparison or because they are crucial inputs in the proofs of Strichartz inequalities
- Present some recent results (joint with H. Mizutani) on Strichartz inequalities on asymptotically flat manifolds

(ロ)、(型)、(E)、(E)、 E) の(の)

Strichartz inequalities for the Schrödinger equation

Strichartz inequalities for the Schrödinger equation take the form

$$\left(\int_{-T}^{T} ||e^{it\Delta}u_0||_{L^q(\mathbb{R}^n)}^p dt\right)^{\frac{1}{p}} \leq C||u_0||_{L^2}$$

◆□ > ◆□ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ●

provided (p, q) is **admissible** (scaling condition)

$$\frac{2}{p} + \frac{n}{q} = \frac{n}{2},$$

Strichartz inequalities for the Schrödinger equation take the form

$$\left(\int_{-\tau}^{\tau} ||e^{it\Delta}u_0||_{L^q(\mathbb{R}^n)}^p dt\right)^{\frac{1}{p}} \leq C||u_0||_{L^2}$$

provided (p, q) is **admissible** (scaling condition)

$$\frac{2}{p} + \frac{n}{q} = \frac{n}{2}, \qquad p, q \ge 2, \ q \neq \infty \text{ if } n = 2.$$

[Strichartz], [Ginibre-Velo], [Keel-Tao]

Strichartz inequalities for the Schrödinger equation take the form

$$\left(\int_{-\tau}^{\tau} ||e^{it\Delta}u_0||_{L^q(\mathbb{R}^n)}^p dt\right)^{\frac{1}{p}} \leq C||u_0||_{L^2}$$

provided (p, q) is admissible (scaling condition)

$$\frac{2}{p} + \frac{n}{q} = \frac{n}{2}, \qquad p, q \ge 2, \ q \neq \infty \text{ if } n = 2.$$

[Strichartz], [Ginibre-Velo], [Keel-Tao] Interests:

1. Shows that $e^{it\Delta}u_0 \in L^q$ for a.e. t without using any derivative on u_0 .

Strichartz inequalities for the Schrödinger equation take the form

$$\left(\int_{-\tau}^{\tau} ||e^{it\Delta}u_0||_{L^q(\mathbb{R}^n)}^p dt\right)^{\frac{1}{p}} \leq C||u_0||_{L^2}$$

provided (p, q) is admissible (scaling condition)

$$\frac{2}{p} + \frac{n}{q} = \frac{n}{2}, \qquad p, q \ge 2, \ q \neq \infty \text{ if } n = 2.$$

[Strichartz], [Ginibre-Velo], [Keel-Tao] Interests:

1. Shows that $e^{it\Delta}u_0 \in L^q$ for a.e. t without using any derivative on u_0 . Compare with Sobolev inequalities

Strichartz inequalities for the Schrödinger equation take the form

$$\left(\int_{-\tau}^{\tau} ||e^{it\Delta}u_0||_{L^q(\mathbb{R}^n)}^p dt\right)^{\frac{1}{p}} \leq C||u_0||_{L^2}$$

provided (p, q) is admissible (scaling condition)

$$\frac{2}{p} + \frac{n}{q} = \frac{n}{2}, \qquad p, q \ge 2, \ q \neq \infty \text{ if } n = 2.$$

[Strichartz], [Ginibre-Velo], [Keel-Tao] Interests:

1. Shows that $e^{it\Delta}u_0 \in L^q$ for a.e. t without using any derivative on u_0 . Compare with Sobolev inequalities $(2 \le q < \infty)$

$$||e^{it\Delta}u_0||_{L^q} \lesssim ||e^{it\Delta}u_0||_{H^s} = ||u_0||_{H^s}, \qquad s = \frac{n}{2} - \frac{n}{q}$$

Strichartz inequalities for the Schrödinger equation take the form

$$\left(\int_{-\tau}^{\tau} ||e^{it\Delta}u_0||_{L^q(\mathbb{R}^n)}^p dt\right)^{\frac{1}{p}} \leq C||u_0||_{L^2}$$

provided (p, q) is admissible (scaling condition)

$$\frac{2}{p} + \frac{n}{q} = \frac{n}{2}, \qquad p, q \ge 2, \ q \neq \infty \text{ if } n = 2.$$

[Strichartz], [Ginibre-Velo], [Keel-Tao] Interests:

1. Shows that $e^{it\Delta}u_0 \in L^q$ for a.e. t without using any derivative on u_0 . Compare with Sobolev inequalities $(2 \le q < \infty)$

$$||e^{it\Delta}u_0||_{L^q} \lesssim ||e^{it\Delta}u_0||_{H^s} = ||u_0||_{H^s}, \qquad s = \frac{n}{2} - \frac{n}{q}$$

2. Important to solve non linear equations at low regularity

Strichartz inequalities for the Schrödinger equation take the form

$$\left(\int_{-\tau}^{\tau} ||e^{it\Delta}u_0||_{L^q(\mathbb{R}^n)}^p dt\right)^{\frac{1}{p}} \leq C||u_0||_{L^2}$$

provided (p, q) is admissible (scaling condition)

$$\frac{2}{p} + \frac{n}{q} = \frac{n}{2}, \qquad p, q \ge 2, \ q \neq \infty \text{ if } n = 2.$$

[Strichartz], [Ginibre-Velo], [Keel-Tao] Interests:

1. Shows that $e^{it\Delta}u_0 \in L^q$ for a.e. t without using any derivative on u_0 . Compare with Sobolev inequalities $(2 \le q < \infty)$

$$||e^{it\Delta}u_0||_{L^q} \lesssim ||e^{it\Delta}u_0||_{H^s} = ||u_0||_{H^s}, \qquad s = \frac{n}{2} - \frac{n}{q}$$

< ロ > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □

- 2. Important to solve non linear equations at low regularity
- 3. For $T = +\infty$

Strichartz inequalities for the Schrödinger equation take the form

$$\left(\int_{-\tau}^{\tau} ||e^{it\Delta}u_0||_{L^q(\mathbb{R}^n)}^p dt\right)^{\frac{1}{p}} \leq C||u_0||_{L^2}$$

provided (p, q) is admissible (scaling condition)

$$\frac{2}{p} + \frac{n}{q} = \frac{n}{2}, \qquad p, q \ge 2, \ q \neq \infty \text{ if } n = 2.$$

[Strichartz], [Ginibre-Velo], [Keel-Tao] Interests:

1. Shows that $e^{it\Delta}u_0 \in L^q$ for a.e. t without using any derivative on u_0 . Compare with Sobolev inequalities $(2 \le q < \infty)$

$$||e^{it\Delta}u_0||_{L^q} \lesssim ||e^{it\Delta}u_0||_{H^s} = ||u_0||_{H^s}, \qquad s = \frac{n}{2} - \frac{n}{q}$$

- 2. Important to solve non linear equations at low regularity
- 3. For $T = +\infty$ (= global in time estimates),

Strichartz inequalities for the Schrödinger equation take the form

$$\left(\int_{-\tau}^{\tau} ||e^{it\Delta}u_0||_{L^q(\mathbb{R}^n)}^p dt\right)^{\frac{1}{p}} \leq C||u_0||_{L^2}$$

provided (p, q) is admissible (scaling condition)

$$\frac{2}{p} + \frac{n}{q} = \frac{n}{2}, \qquad p, q \ge 2, \ q \neq \infty \text{ if } n = 2.$$

[Strichartz], [Ginibre-Velo], [Keel-Tao] Interests:

1. Shows that $e^{it\Delta}u_0 \in L^q$ for a.e. t without using any derivative on u_0 . Compare with Sobolev inequalities $(2 \le q < \infty)$

$$||e^{it\Delta}u_0||_{L^q} \lesssim ||e^{it\Delta}u_0||_{H^s} = ||u_0||_{H^s}, \qquad s = \frac{n}{2} - \frac{n}{q}$$

- 2. Important to solve non linear equations at low regularity
- 3. For $T = +\infty$ (= global in time estimates), shows that $||e^{it\Delta}u_0||_{L^q} \to 0$ as $t \to \infty$ (on L^{ρ} average if q > 2)

Strichartz inequalities for the Schrödinger equation take the form

$$\left(\int_{-\tau}^{\tau} ||e^{it\Delta}u_0||_{L^q(\mathbb{R}^n)}^p dt\right)^{\frac{1}{p}} \leq C||u_0||_{L^2}$$

provided (p, q) is admissible (scaling condition)

$$\frac{2}{p} + \frac{n}{q} = \frac{n}{2}, \qquad p, q \ge 2, \ q \neq \infty \text{ if } n = 2.$$

[Strichartz], [Ginibre-Velo], [Keel-Tao] Interests:

1. Shows that $e^{it\Delta}u_0 \in L^q$ for a.e. t without using any derivative on u_0 . Compare with Sobolev inequalities $(2 \le q < \infty)$

$$||e^{it\Delta}u_0||_{L^q} \lesssim ||e^{it\Delta}u_0||_{H^s} = ||u_0||_{H^s}, \qquad s = \frac{n}{2} - \frac{n}{q}$$

- 2. Important to solve non linear equations at low regularity
- 3. For $T = +\infty$ (= global in time estimates), shows that $||e^{it\Delta}u_0||_{L^q} \to 0$ as $t \to \infty$ (on L^p average if q > 2) ~ local energy decay (RAGE Theorem)

Strichartz inequalities for the Schrödinger equation take the form

$$\left(\int_{-\tau}^{\tau} ||e^{it\Delta}u_0||_{L^q(\mathbb{R}^n)}^p dt\right)^{\frac{1}{p}} \leq C||u_0||_{L^2}$$

provided (p, q) is admissible (scaling condition)

$$\frac{2}{p} + \frac{n}{q} = \frac{n}{2}, \qquad p, q \ge 2, \ q \neq \infty \text{ if } n = 2.$$

[Strichartz], [Ginibre-Velo], [Keel-Tao] Interests:

1. Shows that $e^{it\Delta}u_0 \in L^q$ for a.e. t without using any derivative on u_0 . Compare with Sobolev inequalities $(2 \le q < \infty)$

$$||e^{it\Delta}u_0||_{L^q} \lesssim ||e^{it\Delta}u_0||_{H^s} = ||u_0||_{H^s}, \qquad s = \frac{n}{2} - \frac{n}{q}$$

- 2. Important to solve non linear equations at low regularity
- 3. For $T = +\infty$ (= global in time estimates), shows that $||e^{it\Delta}u_0||_{L^q} \to 0$ as $t \to \infty$ (on L^p average if q > 2) ~ local energy decay (RAGE Theorem) since

$$||e^{it\Delta}u_0||_{L^2(K)} \lesssim_K ||e^{it\Delta}u_0||_{L^q(\mathbb{R}^n)}, \qquad K \in \mathbb{R}^n$$

Scattering inequalities

(ロ)、(型)、(E)、(E)、 E) の(の)

Scattering inequalities

Resolvent estimates:

Scattering inequalities

 \blacktriangleright Resolvent estimates: give the behaviour with respect to $\lambda \in \mathbb{R}$ of

$$R_0(\lambda \pm i0) = \lim_{\delta \to 0^{\pm}} (-\Delta - \lambda - i\delta)^{-1}$$

(ロ)、(型)、(E)、(E)、 E) の(の)

Scattering inequalities

 \blacktriangleright Resolvent estimates: give the behaviour with respect to $\lambda \in \mathbb{R}$ of

$$R_0(\lambda \pm i0) = \lim_{\delta \to 0^{\pm}} (-\Delta - \lambda - i\delta)^{-1}$$

(ロ)、(型)、(E)、(E)、 E) の(の)

In general, the existence of the limit is called limiting absorption principle

Scattering inequalities

▶ Resolvent estimates: give the behaviour with respect to $\lambda \in \mathbb{R}$ of

$$R_0(\lambda \pm i0) = \lim_{\delta \to 0^{\pm}} (-\Delta - \lambda - i\delta)^{-1}$$

In general, the existence of the limit is called **limiting absorption principle** Intuition. $R_0(\lambda + i\delta)$ is the Fourier multiplier by $(|\xi|^2 - \lambda - i\delta)^{-1}$.

Scattering inequalities

 \blacktriangleright Resolvent estimates: give the behaviour with respect to $\lambda \in \mathbb{R}$ of

$$R_0(\lambda \pm i0) = \lim_{\delta \to 0^{\pm}} (-\Delta - \lambda - i\delta)^{-1}$$

In general, the existence of the limit is called **limiting absorption principle** <u>Intuition</u>. $R_0(\lambda + i\delta)$ is the Fourier multiplier by $(|\xi|^2 - \lambda - i\delta)^{-1}$. This multiplier has a limit as $\delta \to 0^{\pm}$ (~ principal value) provided it is tested against smooth enough functions

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

Scattering inequalities

▶ Resolvent estimates: give the behaviour with respect to $\lambda \in \mathbb{R}$ of

$$R_0(\lambda \pm i0) = \lim_{\delta \to 0^{\pm}} (-\Delta - \lambda - i\delta)^{-1}$$

In general, the existence of the limit is called **limiting absorption principle** <u>Intuition</u>. $R_0(\lambda + i\delta)$ is the Fourier multiplier by $(|\xi|^2 - \lambda - i\delta)^{-1}$. This multiplier has a limit as $\delta \to 0^{\pm}$ (~ principal value) provided it is tested against smooth enough functions on the Fourier side

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

Scattering inequalities

 \blacktriangleright Resolvent estimates: give the behaviour with respect to $\lambda \in \mathbb{R}$ of

$$R_0(\lambda \pm i0) = \lim_{\delta \to 0^{\pm}} (-\Delta - \lambda - i\delta)^{-1}$$

In general, the existence of the limit is called **limiting absorption principle** <u>Intuition</u>. $R_0(\lambda + i\delta)$ is the Fourier multiplier by $(|\xi|^2 - \lambda - i\delta)^{-1}$. This multiplier has a limit as $\delta \to 0^{\pm}$ (~ principal value) provided it is tested against smooth enough functions on the Fourier side \leftrightarrow decaying functions on the spatial side.

Scattering inequalities

▶ Resolvent estimates: give the behaviour with respect to $\lambda \in \mathbb{R}$ of

$$R_0(\lambda \pm i0) = \lim_{\delta \to 0^{\pm}} (-\Delta - \lambda - i\delta)^{-1}$$

In general, the existence of the limit is called **limiting absorption principle** <u>Intuition</u>. $R_0(\lambda + i\delta)$ is the Fourier multiplier by $(|\xi|^2 - \lambda - i\delta)^{-1}$. This multiplier has a limit as $\delta \to 0^{\pm}$ (~ principal value) provided it is tested against smooth enough functions on the Fourier side \leftrightarrow decaying functions on the spatial side. Examples.

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

1. High energy estimates:

Scattering inequalities

 \blacktriangleright Resolvent estimates: give the behaviour with respect to $\lambda \in \mathbb{R}$ of

$$R_0(\lambda \pm i0) = \lim_{\delta \to 0^{\pm}} (-\Delta - \lambda - i\delta)^{-1}$$

In general, the existence of the limit is called **limiting absorption principle** <u>Intuition</u>. $R_0(\lambda + i\delta)$ is the Fourier multiplier by $(|\xi|^2 - \lambda - i\delta)^{-1}$. This multiplier has a limit as $\delta \to 0^{\pm}$ (~ principal value) provided it is tested against smooth enough functions on the Fourier side \leftrightarrow decaying functions on the spatial side. Examples.

1. High energy estimates: if $\nu > 1/2$,

$$\left|\left|\langle x\rangle^{-\nu}R_0(\lambda\pm i0)\langle x\rangle^{-\nu}\right|\right|_{L^2\to L^2}\lesssim \lambda^{-1/2}, \ \lambda\geq 1$$

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

Scattering inequalities

▶ Resolvent estimates: give the behaviour with respect to $\lambda \in \mathbb{R}$ of

$$R_0(\lambda \pm i0) = \lim_{\delta \to 0^{\pm}} (-\Delta - \lambda - i\delta)^{-1}$$

In general, the existence of the limit is called **limiting absorption principle** <u>Intuition</u>. $R_0(\lambda + i\delta)$ is the Fourier multiplier by $(|\xi|^2 - \lambda - i\delta)^{-1}$. This multiplier has a limit as $\delta \to 0^{\pm}$ (~ principal value) provided it is tested against smooth enough functions on the Fourier side \leftrightarrow decaying functions on the spatial side. Examples.

1. High energy estimates: if $\nu > 1/2$,

$$\left|\left|\langle x\rangle^{-\nu}R_0(\lambda\pm i0)\langle x\rangle^{-\nu}\right|\right|_{L^2\to L^2}\lesssim \lambda^{-1/2}, \ \lambda\geq 1$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ □ □ のへで

2. Low energy estimates:

Scattering inequalities

• Resolvent estimates: give the behaviour with respect to $\lambda \in \mathbb{R}$ of

$$R_0(\lambda \pm i0) = \lim_{\delta \to 0^{\pm}} (-\Delta - \lambda - i\delta)^{-1}$$

In general, the existence of the limit is called **limiting absorption principle** <u>Intuition</u>. $R_0(\lambda + i\delta)$ is the Fourier multiplier by $(|\xi|^2 - \lambda - i\delta)^{-1}$. This multiplier has a limit as $\delta \to 0^{\pm}$ (~ principal value) provided it is tested against smooth enough functions on the Fourier side \leftrightarrow decaying functions on the spatial side. Examples.

1. High energy estimates: if $\nu > 1/2$,

$$\left|\left|\langle x\rangle^{-\nu}R_0(\lambda\pm i0)\langle x\rangle^{-\nu}\right|\right|_{L^2\to L^2}\lesssim \lambda^{-1/2}, \ \lambda\geq 1$$

2. Low energy estimates: if $\nu = 1$ and $n \ge 3$

 $\left|\left|\langle x\rangle^{-1}R_0(\lambda\pm i0)\langle x\rangle^{-1}\right|\right|_{L^2\to L^2}\lesssim 1, \ |\lambda|\leq 1$

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

Scattering inequalities

▶ Resolvent estimates: give the behaviour with respect to $\lambda \in \mathbb{R}$ of

$$R_0(\lambda \pm i0) = \lim_{\delta \to 0^{\pm}} (-\Delta - \lambda - i\delta)^{-1}$$

In general, the existence of the limit is called **limiting absorption principle** <u>Intuition</u>. $R_0(\lambda + i\delta)$ is the Fourier multiplier by $(|\xi|^2 - \lambda - i\delta)^{-1}$. This multiplier has a limit as $\delta \to 0^{\pm}$ (~ principal value) provided it is tested against smooth enough functions on the Fourier side \leftrightarrow decaying functions on the spatial side. Examples.

1. High energy estimates: if $\nu > 1/2$,

$$\left|\left|\langle x\rangle^{-\nu}R_0(\lambda\pm i0)\langle x\rangle^{-\nu}\right|\right|_{L^2\to L^2}\lesssim \lambda^{-1/2}, \ \lambda\geq 1$$

2. Low energy estimates: if $\nu = 1$ and $n \ge 3$

$$\left|\left|\langle x\rangle^{-1}R_0(\lambda\pm i0)\langle x\rangle^{-1}\right|\right|_{L^2\to L^2}\lesssim 1, \ |\lambda|\leq 1$$

3. One may (actually, one has to) also consider estimates on

$$R_0(\lambda \pm i0)^{\prime}$$

Scattering inequalities

▶ Resolvent estimates: give the behaviour with respect to $\lambda \in \mathbb{R}$ of

$$R_0(\lambda \pm i0) = \lim_{\delta \to 0^{\pm}} (-\Delta - \lambda - i\delta)^{-1}$$

In general, the existence of the limit is called **limiting absorption principle** <u>Intuition</u>. $R_0(\lambda + i\delta)$ is the Fourier multiplier by $(|\xi|^2 - \lambda - i\delta)^{-1}$. This multiplier has a limit as $\delta \to 0^{\pm}$ (~ principal value) provided it is tested against smooth enough functions on the Fourier side \leftrightarrow decaying functions on the spatial side. Examples.

1. High energy estimates: if $\nu > 1/2$,

$$\left|\left|\langle x\rangle^{-\nu}R_0(\lambda\pm i0)\langle x\rangle^{-\nu}\right|\right|_{L^2\to L^2}\lesssim \lambda^{-1/2}, \ \lambda\geq 1$$

2. Low energy estimates: if $\nu = 1$ and $n \ge 3$

$$\left|\left|\langle x\rangle^{-1}R_0(\lambda\pm i0)\langle x\rangle^{-1}\right|\right|_{L^2\to L^2}\lesssim 1, \ |\lambda|\leq 1$$

3. One may (actually, one has to) also consider estimates on

$$R_0(\lambda \pm i0)^k = \frac{1}{(k-1)!} \frac{d^{k-1}}{d\lambda^{k-1}} R_0(\lambda \pm i0)$$

(ロ)、(型)、(E)、(E)、 E) の(の)

Scattering inequalities (continued)

Propagation / time decay estimates:

Scattering inequalities (continued)

▶ Propagation / time decay estimates: given a (spectral) cutoff $\varphi \in C_0^{\infty}(0, +\infty)$, understand the time decay of

$$\varphi(-\Delta/\lambda)e^{it\Delta}$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

Scattering inequalities (continued)

▶ Propagation / time decay estimates: given a (spectral) cutoff $\varphi \in C_0^{\infty}(0, +\infty)$, understand the time decay of

$$\varphi(-\Delta/\lambda)e^{it\Delta}$$

as $t \to \infty$,

Scattering inequalities (continued)

▶ Propagation / time decay estimates: given a (spectral) cutoff $\varphi \in C_0^{\infty}(0, +\infty)$, understand the time decay of

$$arphi(-\Delta/\lambda)e^{it\Delta}$$

as $t \to \infty$, in term of the parameter $\lambda > 0$.

Scattering inequalities (continued)

▶ Propagation / time decay estimates: given a (spectral) cutoff $\varphi \in C_0^{\infty}(0, +\infty)$, understand the time decay of

$$arphi(-\Delta/\lambda)e^{it\Delta}$$

as $t \to \infty$, in term of the parameter $\lambda > 0$. Intuition. For $\lambda = 1$, the Schwartz kernel of $\varphi(-\Delta)e^{it\Delta}$ is the oscillatory integral

▲□▶ ▲□▶ ▲□▶ ▲□▶ □ □ のへで

$$\int e^{i(x-y)\cdot\xi-it|\xi|^2}\varphi(|\xi|^2)\frac{d\xi}{(2\pi)^n}$$

Scattering inequalities (continued)

▶ Propagation / time decay estimates: given a (spectral) cutoff $\varphi \in C_0^{\infty}(0, +\infty)$, understand the time decay of

$$arphi(-\Delta/\lambda)e^{it\Delta}$$

as $t \to \infty$, in term of the parameter $\lambda > 0$. Intuition. For $\lambda = 1$, the Schwartz kernel of $\varphi(-\Delta)e^{it\Delta}$ is the oscillatory integral

$$\int e^{i(x-y)\cdot\xi-it|\xi|^2}\varphi(|\xi|^2)\frac{d\xi}{(2\pi)^n} = \frac{i}{2t}\int \left(\frac{\xi}{2|\xi|^2}\cdot\partial_\xi e^{-it|\xi|^2}\right)e^{i(x-y)\cdot\xi}\varphi(|\xi|^2)\frac{d\xi}{(2\pi)^n}$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ □ □ のへで

Scattering inequalities (continued)

▶ Propagation / time decay estimates: given a (spectral) cutoff $\varphi \in C_0^\infty(0, +\infty)$, understand the time decay of

$$arphi(-\Delta/\lambda)e^{it\Delta}$$

as $t \to \infty$, in term of the parameter $\lambda > 0$. Intuition. For $\lambda = 1$, the Schwartz kernel of $\varphi(-\Delta)e^{it\Delta}$ is the oscillatory integral

$$\int e^{i(x-y)\cdot\xi-it|\xi|^2}\varphi(|\xi|^2)\frac{d\xi}{(2\pi)^n} = \frac{i}{2t}\int \left(\frac{\xi}{2|\xi|^2}\cdot\partial_\xi e^{-it|\xi|^2}\right)e^{i(x-y)\cdot\xi}\varphi(|\xi|^2)\frac{d\xi}{(2\pi)^n}$$

which leads to

$$\left|\left|\langle x\rangle^{-k}\varphi(-\Delta)e^{it\Delta}\langle x\rangle^{-k}\right|\right|_{L^2\to L^2}\lesssim \langle t\rangle^{-k}$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ □ □ のへで

Scattering inequalities (continued)

▶ Propagation / time decay estimates: given a (spectral) cutoff $\varphi \in C_0^\infty(0, +\infty)$, understand the time decay of

$$arphi(-\Delta/\lambda)e^{it\Delta}$$

as $t \to \infty$, in term of the parameter $\lambda > 0$. Intuition. For $\lambda = 1$, the Schwartz kernel of $\varphi(-\Delta)e^{it\Delta}$ is the oscillatory integral

$$\int e^{i(x-y)\cdot\xi-it|\xi|^2}\varphi(|\xi|^2)\frac{d\xi}{(2\pi)^n} = \frac{i}{2t}\int \left(\frac{\xi}{2|\xi|^2}\cdot\partial_{\xi}e^{-it|\xi|^2}\right)e^{i(x-y)\cdot\xi}\varphi(|\xi|^2)\frac{d\xi}{(2\pi)^n}$$

which leads to

$$\left|\left|\langle x\rangle^{-k}\varphi(-\Delta)e^{it\Delta}\langle x\rangle^{-k}\right|\right|_{L^2\to L^2}\lesssim \langle t\rangle^{-k}.$$

By scaling

$$\left|\left|\langle\lambda^{\frac{1}{2}}x\rangle^{-k}\varphi(-\Delta/\lambda)e^{it\Delta}\langle\lambda^{\frac{1}{2}}x\rangle^{-k}\right|\right|_{L^{2}\to L^{2}}\lesssim\langle\lambda t\rangle^{-k}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへ⊙
Scattering inequalities (end)

Integrated decay/ smoothing estimates:

Scattering inequalities (end)

Integrated decay/ smoothing estimates: Integrated space-time decay estimates are of the form

$$\left(\int_{\mathbb{R}} ||\langle x\rangle^{-\nu}\varphi(-\Delta/\lambda)e^{it\Delta}u_0||_{L^2}^2 dt\right)^{\frac{1}{2}} \lesssim_{\lambda} ||u_0||_{L^2},$$

(ロ)、(型)、(E)、(E)、 E) の(の)

with $\nu > 1/2$.

Scattering inequalities (end)

 Integrated decay/ smoothing estimates: Integrated space-time decay estimates are of the form

$$\left(\int_{\mathbb{R}} ||\langle x\rangle^{-\nu}\varphi(-\Delta/\lambda)e^{it\Delta}u_0||_{L^2}^2 dt\right)^{\frac{1}{2}} \lesssim_{\lambda} ||u_0||_{L^2},$$

with $\nu > 1/2$. By tracking the dependence on λ , one may obtain the non spectrally localized estimate ($n \ge 3$)

$$\left(\int_{\mathbb{R}} ||\langle x\rangle^{-1} \langle D\rangle^{\frac{1}{2}} e^{it\Delta} u_0||_{L^2}^2 dt\right)^{\frac{1}{2}} \lesssim ||u_0||_{L^2}$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ □臣 = のへで

Scattering inequalities (end)

 Integrated decay/ smoothing estimates: Integrated space-time decay estimates are of the form

$$\left(\int_{\mathbb{R}} ||\langle x\rangle^{-\nu} \varphi(-\Delta/\lambda) e^{it\Delta} u_0||_{L^2}^2 dt\right)^{\frac{1}{2}} \lesssim_{\lambda} ||u_0||_{L^2},$$

with $\nu > 1/2$. By tracking the dependence on λ , one may obtain the non spectrally localized estimate ($n \ge 3$)

$$\left(\int_{\mathbb{R}} ||\langle x\rangle^{-1} \langle D\rangle^{\frac{1}{2}} e^{it\Delta} u_0||_{L^2}^2 dt\right)^{\frac{1}{2}} \lesssim ||u_0||_{L^2}$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

which is the $\frac{1}{2}$ -smoothing effect for the Schrödinger equation.

Scattering inequalities (end)

 Integrated decay/ smoothing estimates: Integrated space-time decay estimates are of the form

$$\left(\int_{\mathbb{R}} ||\langle x\rangle^{-\nu}\varphi(-\Delta/\lambda)e^{it\Delta}u_0||_{L^2}^2 dt\right)^{\frac{1}{2}} \lesssim_{\lambda} ||u_0||_{L^2},$$

with $\nu > 1/2$. By tracking the dependence on λ , one may obtain the non spectrally localized estimate ($n \ge 3$)

$$\left(\int_{\mathbb{R}} ||\langle x\rangle^{-1} \langle D\rangle^{\frac{1}{2}} e^{it\Delta} u_0||_{L^2}^2 dt\right)^{\frac{1}{2}} \lesssim ||u_0||_{L^2}$$

which is the $\frac{1}{2}$ -smoothing effect for the Schrödinger equation. Note that even locally in time (i.e. with \mathbb{R} replaced by $[-\mathcal{T}, \mathcal{T}]$) this is non trivial.

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

Scattering inequalities (end)

 Integrated decay/ smoothing estimates: Integrated space-time decay estimates are of the form

$$\left(\int_{\mathbb{R}} ||\langle x\rangle^{-\nu}\varphi(-\Delta/\lambda)e^{it\Delta}u_0||_{L^2}^2 dt\right)^{\frac{1}{2}} \lesssim_{\lambda} ||u_0||_{L^2},$$

with $\nu > 1/2$. By tracking the dependence on λ , one may obtain the non spectrally localized estimate ($n \ge 3$)

$$\left(\int_{\mathbb{R}} ||\langle x\rangle^{-1} \langle D\rangle^{\frac{1}{2}} e^{it\Delta} u_0||_{L^2}^2 dt\right)^{\frac{1}{2}} \lesssim ||u_0||_{L^2}$$

which is the $\frac{1}{2}$ -smoothing effect for the Schrödinger equation. Note that even locally in time (i.e. with \mathbb{R} replaced by $[-\mathcal{T}, \mathcal{T}]$) this is non trivial.

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

Intuition. More on the next slides.

Scattering inequalities (end)

 Integrated decay/ smoothing estimates: Integrated space-time decay estimates are of the form

$$\left(\int_{\mathbb{R}} ||\langle x\rangle^{-\nu}\varphi(-\Delta/\lambda)e^{it\Delta}u_0||_{L^2}^2 dt\right)^{\frac{1}{2}} \lesssim_{\lambda} ||u_0||_{L^2},$$

with $\nu > 1/2$. By tracking the dependence on λ , one may obtain the non spectrally localized estimate ($n \ge 3$)

$$\left(\int_{\mathbb{R}} ||\langle x\rangle^{-1} \langle D\rangle^{\frac{1}{2}} e^{it\Delta} u_0||_{L^2}^2 dt\right)^{\frac{1}{2}} \lesssim ||u_0||_{L^2}$$

which is the $\frac{1}{2}$ -smoothing effect for the Schrödinger equation. Note that even locally in time (i.e. with \mathbb{R} replaced by $[-\tau, \tau]$) this is non trivial.

Intuition. More on the next slides. Technically, they follow from resolvent estimates via a Parseval argument, using that $e^{it\Delta}$ is the Fourier transform $(\lambda \to t)$ of the spectral measure

$$R_0(\lambda + i0) - R_0(\lambda - i0).$$

Scattering inequalities (end)

 Integrated decay/ smoothing estimates: Integrated space-time decay estimates are of the form

$$\left(\int_{\mathbb{R}} ||\langle x\rangle^{-\nu}\varphi(-\Delta/\lambda)e^{it\Delta}u_0||_{L^2}^2 dt\right)^{\frac{1}{2}} \lesssim_{\lambda} ||u_0||_{L^2},$$

with $\nu > 1/2$. By tracking the dependence on λ , one may obtain the non spectrally localized estimate ($n \geq 3$)

$$\left(\int_{\mathbb{R}} ||\langle x\rangle^{-1} \langle D\rangle^{\frac{1}{2}} e^{it\Delta} u_0||_{L^2}^2 dt\right)^{\frac{1}{2}} \lesssim ||u_0||_{L^2}$$

which is the $\frac{1}{2}$ -smoothing effect for the Schrödinger equation. Note that even locally in time (i.e. with \mathbb{R} replaced by $[-\tau, \tau]$) this is non trivial.

Intuition. More on the next slides. Technically, they follow from resolvent estimates via a Parseval argument, using that $e^{it\Delta}$ is the Fourier transform $(\lambda \to t)$ of the spectral measure

$$R_0(\lambda + i0) - R_0(\lambda - i0).$$

<u>Rem</u>. This correspondence $\lambda \to t$ also allows to convert resolvent estimates into time decay/propagation estimates (smoothness of $R_0(\lambda \pm i0) \leftrightarrow$ decay of e^{itP})

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

<□ > < @ > < E > < E > E のQ @

Strichartz inequalities

Strichartz inequalities

Consider the L^2 normalized semiclassical wave packet

$$G_{z,\zeta,h}(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x-z) - \frac{|x-z|^2}{2h}\right)$$

Strichartz inequalities

Consider the L^2 normalized semiclassical wave packet

$$G_{z,\zeta,h}(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x-z) - \frac{|x-z|^2}{2h}\right)$$

(ロ)、(型)、(E)、(E)、 E) の(の)

Then,

$$\left|e^{i\frac{t}{2}\Delta}G_{z,\zeta,h}(x)\right|$$

Strichartz inequalities

Consider the L^2 normalized semiclassical wave packet

$$G_{z,\zeta,h}(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x-z) - \frac{|x-z|^2}{2h}\right)$$

Then,

$$\left|e^{i\frac{t}{2}\Delta}G_{z,\zeta,h}(x)\right| = \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h\rangle^{2}\right)^{\frac{n}{4}}}\exp\left(-\frac{\left|x-z-(t/h)\zeta\right|^{2}}{2h\langle t/h\rangle^{2}}\right)$$

(ロ)、(型)、(E)、(E)、 E) の(の)

with $\langle \tau \rangle = (1 + \tau^2)^{\frac{1}{2}}$.

Strichartz inequalities

Consider the L^2 normalized semiclassical wave packet

$$G_{z,\zeta,h}(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x-z) - \frac{|x-z|^2}{2h}\right)$$

Then,

$$\left|e^{j\frac{t}{2}\Delta}G_{z,\zeta,h}(x)\right| = \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h\rangle^{2}\right)^{\frac{n}{4}}}\exp\left(-\frac{\left|x-z-(t/h)\zeta\right|^{2}}{2h\langle t/h\rangle^{2}}\right)$$

with $\langle au
angle = (1+ au^2)^{rac{1}{2}}.$ This implies easily

$$\left|\left|e^{i\frac{t}{2}\Delta}G_{z,\zeta,h}\right|\right|_{L^{q}} = (2/q)^{\frac{n}{2q}} \left(\frac{1}{\pi h \langle t/h \rangle^{2}}\right)^{\frac{n}{2}\left(\frac{1}{2} - \frac{1}{q}\right)}$$

◆□ > ◆□ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ◆ □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ● □ > ●

Strichartz inequalities

Consider the L^2 normalized semiclassical wave packet

$$G_{z,\zeta,h}(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x-z) - \frac{|x-z|^2}{2h}\right)$$

Then,

$$\left|e^{j\frac{t}{2}\Delta}G_{z,\zeta,h}(x)\right| = \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h\rangle^2\right)^{\frac{n}{4}}}\exp\left(-\frac{\left|x-z-(t/h)\zeta\right|^2}{2h\langle t/h\rangle^2}\right)$$

with $\langle au
angle = (1+ au^2)^{rac{1}{2}}.$ This implies easily

$$\left|\left|e^{i\frac{t}{2}\Delta}G_{z,\zeta,h}\right|\right|_{L^{q}} = (2/q)^{\frac{n}{2q}} \left(\frac{1}{\pi h \langle t/h \rangle^{2}}\right)^{\frac{n}{2}\left(\frac{1}{2} - \frac{1}{q}\right)}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへで

Remark. The translation by $(t/h)\zeta$ is **not** used. Only the spreading/dilation factor $\langle t/h \rangle$ plays a role.

Strichartz inequalities

Consider the L^2 normalized semiclassical wave packet

$$G_{z,\zeta,h}(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x-z) - \frac{|x-z|^2}{2h}\right)$$

Then,

$$\left|e^{j\frac{t}{2}\Delta}G_{z,\zeta,h}(x)\right| = \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h\rangle^2\right)^{\frac{n}{4}}}\exp\left(-\frac{\left|x-z-(t/h)\zeta\right|^2}{2h\langle t/h\rangle^2}\right)$$

with $\langle au
angle = (1+ au^2)^{rac{1}{2}}.$ This implies easily

$$\left|\left|e^{i\frac{t}{2}\Delta}G_{z,\zeta,h}\right|\right|_{L^{q}} = (2/q)^{\frac{n}{2q}} \left(\frac{1}{\pi h \langle t/h \rangle^{2}}\right)^{\frac{n}{2}\left(\frac{1}{2} - \frac{1}{q}\right)}$$

(日) (日) (日) (日) (日) (日) (日) (日)

Remark. The translation by $(t/h)\zeta$ is **not** used. Only the spreading/dilation factor $\langle t/h \rangle$ plays a role.

In particular, for $q = 2^* = 2n/(n-2)$,

$$\int_{-T}^{T} \left| \left| e^{i\frac{t}{2}\Delta} G_{z,\zeta,h} \right| \right|_{L^{2*}}^{2} dt$$

Strichartz inequalities

Consider the L^2 normalized semiclassical wave packet

$$G_{z,\zeta,h}(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x-z) - \frac{|x-z|^2}{2h}\right)$$

Then,

$$\left|e^{j\frac{t}{2}\Delta}G_{z,\zeta,h}(x)\right| = \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h\rangle^2\right)^{\frac{n}{4}}}\exp\left(-\frac{\left|x-z-(t/h)\zeta\right|^2}{2h\langle t/h\rangle^2}\right)$$

with $\langle au
angle = (1+ au^2)^{rac{1}{2}}.$ This implies easily

$$\left|\left|e^{i\frac{t}{2}\Delta}G_{z,\zeta,h}\right|\right|_{L^{q}} = (2/q)^{\frac{n}{2q}} \left(\frac{1}{\pi h \langle t/h \rangle^{2}}\right)^{\frac{n}{2}\left(\frac{1}{2} - \frac{1}{q}\right)}$$

(日) (日) (日) (日) (日) (日) (日) (日)

Remark. The translation by $(t/h)\zeta$ is **not** used. Only the spreading/dilation factor $\langle t/h \rangle$ plays a role.

In particular, for $q = 2^* = 2n/(n-2)$,

$$\int_{-\tau}^{\tau} \left| \left| e^{i\frac{t}{2}\Delta} G_{z,\zeta,h} \right| \right|_{L^{2*}}^{2} dt = c_n \int_{-\tau}^{\tau} \frac{1}{\langle t/h \rangle^2} \frac{dt}{h}$$

Strichartz inequalities

Consider the L^2 normalized semiclassical wave packet

$$G_{z,\zeta,h}(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x-z) - \frac{|x-z|^2}{2h}\right)$$

Then,

$$\left|e^{j\frac{t}{2}\Delta}G_{z,\zeta,h}(x)\right| = \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h\rangle^2\right)^{\frac{n}{4}}}\exp\left(-\frac{\left|x-z-(t/h)\zeta\right|^2}{2h\langle t/h\rangle^2}\right)$$

with $\langle au
angle = (1+ au^2)^{rac{1}{2}}.$ This implies easily

$$\left|\left|e^{i\frac{t}{2}\Delta}G_{z,\zeta,h}\right|\right|_{L^{q}} = (2/q)^{\frac{n}{2q}} \left(\frac{1}{\pi h \langle t/h \rangle^{2}}\right)^{\frac{n}{2}\left(\frac{1}{2} - \frac{1}{q}\right)}$$

Remark. The translation by $(t/h)\zeta$ is **not** used. Only the spreading/dilation factor $\langle t/h \rangle$ plays a role.

In particular, for $q = 2^* = 2n/(n-2)$,

$$\int_{-T}^{T} \left| \left| e^{i\frac{t}{2}\Delta} G_{z,\zeta,h} \right| \right|_{L^{2^*}}^2 dt = c_n \int_{-T}^{T} \frac{1}{\langle t/h \rangle^2} \frac{dt}{h} = c_n \int_{-T/h}^{T/h} \frac{1}{1+\tau^2} d\tau \le C$$

Strichartz inequalities

Consider the L^2 normalized semiclassical wave packet

$$G_{z,\zeta,h}(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x-z) - \frac{|x-z|^2}{2h}\right)$$

Then,

$$\left|e^{j\frac{t}{2}\Delta}G_{z,\zeta,h}(x)\right| = \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h\rangle^2\right)^{\frac{n}{4}}}\exp\left(-\frac{\left|x-z-(t/h)\zeta\right|^2}{2h\langle t/h\rangle^2}\right)$$

with $\langle au
angle = (1+ au^2)^{rac{1}{2}}.$ This implies easily

$$\left|\left|e^{i\frac{t}{2}\Delta}G_{z,\zeta,h}\right|\right|_{L^{q}} = (2/q)^{\frac{n}{2q}} \left(\frac{1}{\pi h \langle t/h \rangle^{2}}\right)^{\frac{n}{2}\left(\frac{1}{2} - \frac{1}{q}\right)}$$

Remark. The translation by $(t/h)\zeta$ is **not** used. Only the spreading/dilation factor $\langle t/h \rangle$ plays a role.

In particular, for $q = 2^* = 2n/(n-2)$,

$$\int_{-T}^{T} \left| \left| e^{j\frac{t}{2}\Delta} G_{z,\zeta,h} \right| \right|_{L^{2^*}}^2 dt = c_n \int_{-T}^{T} \frac{1}{\langle t/h \rangle^2} \frac{dt}{h} = c_n \int_{-T/h}^{T/h} \frac{1}{1+\tau^2} d\tau \le C$$

for all $h \in (0, 1]$

Strichartz inequalities

Consider the L^2 normalized semiclassical wave packet

$$G_{z,\zeta,h}(x) = (\pi h)^{-\frac{n}{4}} \exp\left(\frac{i}{h}\zeta \cdot (x-z) - \frac{|x-z|^2}{2h}\right)$$

Then,

$$\left|e^{j\frac{t}{2}\Delta}G_{z,\zeta,h}(x)\right| = \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h\rangle^2\right)^{\frac{n}{4}}}\exp\left(-\frac{\left|x-z-(t/h)\zeta\right|^2}{2h\langle t/h\rangle^2}\right)$$

with $\langle au
angle = (1+ au^2)^{rac{1}{2}}.$ This implies easily

$$\left|\left|e^{i\frac{t}{2}\Delta}G_{z,\zeta,h}\right|\right|_{L^{q}} = (2/q)^{\frac{n}{2q}} \left(\frac{1}{\pi h \langle t/h \rangle^{2}}\right)^{\frac{n}{2}\left(\frac{1}{2} - \frac{1}{q}\right)}$$

Remark. The translation by $(t/h)\zeta$ is **not** used. Only the spreading/dilation factor $\langle t/h \rangle$ plays a role.

In particular, for $q = 2^* = 2n/(n-2)$,

$$\int_{-T}^{T} \left| \left| e^{j\frac{t}{2}\Delta} G_{z,\zeta,h} \right| \right|_{L^{2^*}}^2 dt = c_n \int_{-T}^{T} \frac{1}{\langle t/h \rangle^2} \frac{dt}{h} = c_n \int_{-T/h}^{T/h} \frac{1}{1+\tau^2} d\tau \le C$$

(日) (日) (日) (日) (日) (日) (日) (日)

for all $h \in (0, 1]$ and $z \in \mathbb{R}^n$.

Smoothing effect (local in time)

 $\left|\langle D\rangle^{s}e^{i\frac{t}{2}\Delta}G_{z,\zeta,h}(x)\right|$



Smoothing effect (local in time)

$$\left| \langle D \rangle^s e^{j\frac{t}{2}\Delta} G_{z,\zeta,h}(x) \right| \quad \sim \quad \langle \zeta/h \rangle^s \frac{\pi^{-\frac{n}{4}}}{\left(h \langle t/h \rangle^2\right)^{\frac{n}{4}}} \exp\left(-\frac{\left|x-z-(t/h)\zeta\right|^2}{2h \langle t/h \rangle^2}\right) \qquad h \to 0,$$

Smoothing effect (local in time)

$$\begin{aligned} |\langle D\rangle^{s} e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x)| &\sim \quad \langle \zeta/h \rangle^{s} \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h \rangle^{2}\right)^{\frac{n}{4}}} \exp\left(-\frac{|x-z-(t/h)\zeta|^{2}}{2h\langle t/h \rangle^{2}}\right) \qquad h \to 0, \\ &= \quad \langle \zeta/h \rangle^{s} G_{z,\zeta,h}^{t}(x). \end{aligned}$$

Smoothing effect (local in time)

$$\begin{split} |\langle D \rangle^{s} e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x)| &\sim \quad \langle \zeta/h \rangle^{s} \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h \rangle^{2}\right)^{\frac{n}{4}}} \exp\left(-\frac{|x-z-(t/h)\zeta|^{2}}{2h\langle t/h \rangle^{2}}\right) \qquad h \to 0, \\ &= \quad \langle \zeta/h \rangle^{s} G_{z,\zeta,h}^{t}(x). \end{split}$$

(ロ)、(型)、(E)、(E)、 E) の(の)

We assume that $\zeta \neq 0$,

Smoothing effect (local in time)

$$\begin{split} |\langle D \rangle^{s} e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x)| &\sim \quad \langle \zeta/h \rangle^{s} \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h \rangle^{2}\right)^{\frac{n}{4}}} \exp\left(-\frac{|x-z-(t/h)\zeta|^{2}}{2h\langle t/h \rangle^{2}}\right) \qquad h \to 0, \\ &= \quad \langle \zeta/h \rangle^{s} G_{z,\zeta,h}^{t}(x). \end{split}$$

We assume that $\zeta \neq {\rm 0, \ say} \ |\zeta| = 1$

Smoothing effect (local in time)

$$\begin{aligned} |\langle D \rangle^{s} e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x)| &\sim \quad \langle \zeta/h \rangle^{s} \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h \rangle^{2}\right)^{\frac{n}{4}}} \exp\left(-\frac{|x-z-(t/h)\zeta|^{2}}{2h\langle t/h \rangle^{2}}\right) \qquad h \to 0, \\ &= \quad \langle \zeta/h \rangle^{s} G_{z,\zeta,h}^{t}(x). \end{aligned}$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

We assume that $\zeta\neq 0,$ say $|\zeta|=1$ and then, by possibly rotating the axis, that $\zeta=(1,0,\ldots,0).$

Smoothing effect (local in time)

$$\begin{aligned} |\langle D \rangle^{s} e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x)| &\sim \quad \langle \zeta/h \rangle^{s} \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h \rangle^{2}\right)^{\frac{n}{4}}} \exp\left(-\frac{|x-z-(t/h)\zeta|^{2}}{2h\langle t/h \rangle^{2}}\right) \qquad h \to 0, \\ &= \quad \langle \zeta/h \rangle^{s} G_{z,\zeta,h}^{t}(x). \end{aligned}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへで

We assume that $\zeta\neq 0,$ say $|\zeta|=1$ and then, by possibly rotating the axis, that $\zeta=(1,0,\ldots,0).$ Then

 $\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L^{2}_{x}}^{2}$

Smoothing effect (local in time)

$$\begin{split} |\langle D \rangle^{s} e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x)| &\sim \quad \langle \zeta/h \rangle^{s} \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h \rangle^{2}\right)^{\frac{n}{4}}} \exp\left(-\frac{|x-z-(t/h)\zeta|^{2}}{2h\langle t/h \rangle^{2}}\right) \qquad h \to 0, \\ &= \quad \langle \zeta/h \rangle^{s} G_{z,\zeta,h}^{t}(x). \end{split}$$

We assume that $\zeta\neq 0,$ say $|\zeta|=1$ and then, by possibly rotating the axis, that $\zeta=(1,0,\ldots,0).$ Then

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L_{x}^{2}}^{2}=c_{n}\langle\zeta/h\rangle^{2s}\langle t/h\rangle^{-n}\int\langle h^{\frac{1}{2}}y+z+t\zeta/h\rangle^{-2\nu}\exp\left(-\frac{y^{2}}{\langle t/h\rangle^{2}}\right)dy$$

Smoothing effect (local in time)

$$\begin{split} |\langle D \rangle^{s} e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x)| &\sim \quad \langle \zeta/h \rangle^{s} \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h \rangle^{2}\right)^{\frac{n}{4}}} \exp\left(-\frac{|x-z-(t/h)\zeta|^{2}}{2h\langle t/h \rangle^{2}}\right) \qquad h \to 0, \\ &= \quad \langle \zeta/h \rangle^{s} G_{z,\zeta,h}^{t}(x). \end{split}$$

We assume that $\zeta\neq 0,$ say $|\zeta|=1$ and then, by possibly rotating the axis, that $\zeta=(1,0,\ldots,0).$ Then

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L_{x}^{2}}^{2}=c_{n}\langle\zeta/h\rangle^{2s}\langle t/h\rangle^{-n}\int\langle h^{\frac{1}{2}}y+z+t\zeta/h\rangle^{-2\nu}\exp\left(-\frac{y^{2}}{\langle t/h\rangle^{2}}\right)dy$$

If we further integrate in time on $[-T, T]_t$,

Smoothing effect (local in time)

$$\begin{aligned} |\langle D \rangle^{s} e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x)| &\sim \quad \langle \zeta/h \rangle^{s} \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h \rangle^{2}\right)^{\frac{n}{4}}} \exp\left(-\frac{|x-z-(t/h)\zeta|^{2}}{2h\langle t/h \rangle^{2}}\right) &\quad h \to 0, \\ &= \quad \langle \zeta/h \rangle^{s} G_{z,\zeta,h}^{t}(x). \end{aligned}$$

We assume that $\zeta\neq 0,$ say $|\zeta|=1$ and then, by possibly rotating the axis, that $\zeta=(1,0,\ldots,0).$ Then

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L_{x}^{2}}^{2}=c_{n}\langle\zeta/h\rangle^{2s}\langle t/h\rangle^{-n}\int\langle h^{\frac{1}{2}}y+z+t\zeta/h\rangle^{-2\nu}\exp\left(-\frac{y^{2}}{\langle t/h\rangle^{2}}\right)dy$$

If we further integrate in time on $[-T, T]_t$,

$$c_{n}h\langle\zeta/h\rangle^{2s}\int_{-T/h}^{T/h}\langle\tau\rangle^{-n}\int\langle h^{\frac{1}{2}}y+z+\tau\zeta\rangle^{-2\nu}\exp\left(-\frac{y^{2}}{\langle\tau\rangle^{2}}\right)dyd\tau$$

Smoothing effect (local in time)

$$\begin{aligned} |\langle D \rangle^{s} e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x)| &\sim \quad \langle \zeta/h \rangle^{s} \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h \rangle^{2}\right)^{\frac{n}{4}}} \exp\left(-\frac{|x-z-(t/h)\zeta|^{2}}{2h\langle t/h \rangle^{2}}\right) &\quad h \to 0, \\ &= \quad \langle \zeta/h \rangle^{s} G_{z,\zeta,h}^{t}(x). \end{aligned}$$

We assume that $\zeta\neq 0,$ say $|\zeta|=1$ and then, by possibly rotating the axis, that $\zeta=(1,0,\ldots,0).$ Then

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L_{x}^{2}}^{2}=c_{n}\langle\zeta/h\rangle^{2s}\langle t/h\rangle^{-n}\int\langle h^{\frac{1}{2}}y+z+t\zeta/h\rangle^{-2\nu}\exp\left(-\frac{y^{2}}{\langle t/h\rangle^{2}}\right)dy$$

If we further integrate in time on $[-T, T]_t$,

$$c_n h \langle \zeta/h \rangle^{2s} \int_{-T/h}^{T/h} \langle \tau \rangle^{-n} \int \langle h^{\frac{1}{2}} y + z + \tau \zeta \rangle^{-2\nu} \exp\left(-\frac{y^2}{\langle \tau \rangle^2}\right) dy d\tau$$

which is bounded by

$$c_{n}h\langle 1/h\rangle^{2s}\int_{-T/h}^{T/h}\int \langle h^{\frac{1}{2}}Y_{1}\langle \tau\rangle + z_{1}+\tau\rangle^{-2\nu}\exp\left(-Y^{2}\right)dYd\tau$$

Smoothing effect (local in time)

$$\begin{aligned} |\langle D \rangle^{s} e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x)| &\sim \quad \langle \zeta/h \rangle^{s} \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h \rangle^{2}\right)^{\frac{n}{4}}} \exp\left(-\frac{|x-z-(t/h)\zeta|^{2}}{2h\langle t/h \rangle^{2}}\right) &\quad h \to 0, \\ &= \quad \langle \zeta/h \rangle^{s} G_{z,\zeta,h}^{t}(x). \end{aligned}$$

We assume that $\zeta\neq 0,$ say $|\zeta|=1$ and then, by possibly rotating the axis, that $\zeta=(1,0,\ldots,0).$ Then

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L_{x}^{2}}^{2}=c_{n}\langle\zeta/h\rangle^{2s}\langle t/h\rangle^{-n}\int\langle h^{\frac{1}{2}}y+z+t\zeta/h\rangle^{-2\nu}\exp\left(-\frac{y^{2}}{\langle t/h\rangle^{2}}\right)dy$$

If we further integrate in time on $[-T, T]_t$,

$$c_n h \langle \zeta/h \rangle^{2s} \int_{-T/h}^{T/h} \langle \tau \rangle^{-n} \int \langle h^{\frac{1}{2}} y + z + \tau \zeta \rangle^{-2\nu} \exp\left(-\frac{y^2}{\langle \tau \rangle^2}\right) dy d\tau$$

which is bounded by

$$c_{n}h\langle 1/h\rangle^{2s}\int_{-T/h}^{T/h}\int \langle h^{\frac{1}{2}}Y_{1}\langle \tau\rangle + z_{1}+\tau\rangle^{-2\nu}\exp\left(-Y^{2}\right)dYd\tau$$

Remark. Up to the term $Y_1\langle \tau \rangle$, there is no more contribution of the spreading $\langle \tau \rangle$.

・ロト・西ト・ヨト・ヨー シック

Smoothing effect (local in time)

$$\begin{split} |\langle D \rangle^{s} e^{i\frac{t}{2}\Delta} G_{z,\zeta,h}(x)| &\sim \quad \langle \zeta/h \rangle^{s} \frac{\pi^{-\frac{n}{4}}}{\left(h\langle t/h \rangle^{2}\right)^{\frac{n}{4}}} \exp\left(-\frac{|x-z-(t/h)\zeta|^{2}}{2h\langle t/h \rangle^{2}}\right) \qquad h \to 0, \\ &= \quad \langle \zeta/h \rangle^{s} G_{z,\zeta,h}^{t}(x). \end{split}$$

We assume that $\zeta\neq 0,$ say $|\zeta|=1$ and then, by possibly rotating the axis, that $\zeta=(1,0,\ldots,0).$ Then

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L_{x}^{2}}^{2}=c_{n}\langle\zeta/h\rangle^{2s}\langle t/h\rangle^{-n}\int\langle h^{\frac{1}{2}}y+z+t\zeta/h\rangle^{-2\nu}\exp\left(-\frac{y^{2}}{\langle t/h\rangle^{2}}\right)dy$$

If we further integrate in time on $[-T, T]_t$,

$$c_n h \langle \zeta/h \rangle^{2s} \int_{-T/h}^{T/h} \langle \tau \rangle^{-n} \int \langle h^{\frac{1}{2}} y + z + \tau \zeta \rangle^{-2\nu} \exp\left(-\frac{y^2}{\langle \tau \rangle^2}\right) dy d\tau$$

which is bounded by

$$c_{n}h\langle 1/h\rangle^{2s}\int_{-T/h}^{T/h}\int \langle h^{\frac{1}{2}}Y_{1}\langle \tau\rangle + z_{1}+\tau\rangle^{-2\nu}\exp\left(-Y^{2}\right)dYd\tau$$

Remark. Up to the term $Y_1\langle \tau \rangle$, there is no more contribution of the spreading $\langle \tau \rangle$. Here, the main role will be played the translation by $(t/h)\zeta = \tau \zeta$.

Recall we are estimating

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L^{2}_{t,x}}^{2} \lesssim h^{1-2s}\int_{-T/h}^{T/h}\int\langle h^{\frac{1}{2}}Y_{1}\langle\tau\rangle+z_{1}+\tau\rangle^{-2\nu}\exp\left(-Y^{2}\right)dYd\tau.$$

(ロ)、(型)、(E)、(E)、 E) の(の)

Recall we are estimating

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L^{2}_{t,x}}^{2} \lesssim h^{1-2s}\int_{-T/h}^{T/h}\int\langle h^{\frac{1}{2}}Y_{1}\langle\tau\rangle+z_{1}+\tau\rangle^{-2\nu}\exp\left(-Y^{2}\right)dYd\tau.$$

• In the region $|h^{1/2}Y_1| \leq \epsilon \ (\epsilon \ll 1 \text{ fixed})$,

Recall we are estimating

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L^{2}_{t,x}}^{2} \lesssim h^{1-2s}\int_{-T/h}^{T/h}\int\langle h^{\frac{1}{2}}Y_{1}\langle\tau\rangle+z_{1}+\tau\rangle^{-2\nu}\exp\left(-Y^{2}\right)dYd\tau.$$

▶ In the region $|h^{1/2}Y_1| \le \epsilon$ ($\epsilon \ll 1$ fixed), we integrate in time by using the variable

$$ilde{ au} = au + h^{rac{1}{2}} Y_1 \langle au
angle \qquad ({ t Jacobian} = 1 + O(\epsilon))$$

(ロ)、(型)、(E)、(E)、 E) の(の)

Recall we are estimating

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L^{2}_{t,x}}^{2} \lesssim h^{1-2s}\int_{-T/h}^{T/h}\int\langle h^{\frac{1}{2}}Y_{1}\langle\tau\rangle+z_{1}+\tau\rangle^{-2\nu}\exp\left(-Y^{2}\right)dYd\tau.$$

▶ In the region $|h^{1/2}Y_1| \le \epsilon$ ($\epsilon \ll 1$ fixed), we integrate in time by using the variable

$$ilde{ au} = au + h^{rac{1}{2}} \, Y_1 \langle au
angle \qquad ({ t Jacobian} = 1 + O(\epsilon))$$

so we bound the integral by

$$h^{1-2s} \int \left(\int_{-CT/h}^{CT/h} \langle z_1 + \tilde{\tau} \rangle^{-2\nu} \exp\left(-Y^2\right) d\tilde{\tau} \right) dY \tag{1}$$
Recall we are estimating

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L^{2}_{t,x}}^{2} \lesssim h^{1-2s}\int_{-T/h}^{T/h}\int\langle h^{\frac{1}{2}}Y_{1}\langle\tau\rangle+z_{1}+\tau\rangle^{-2\nu}\exp\left(-Y^{2}\right)dYd\tau.$$

▶ In the region $|h^{1/2}Y_1| \le \epsilon$ ($\epsilon \ll 1$ fixed), we integrate in time by using the variable

$$ilde{ au} = au + h^{rac{1}{2}} \, Y_1 \langle au
angle \qquad ({ t Jacobian} = 1 + O(\epsilon))$$

so we bound the integral by

$$h^{1-2s} \int \left(\int_{-CT/h}^{CT/h} \langle z_1 + \tilde{\tau} \rangle^{-2\nu} \exp\left(-Y^2\right) d\tilde{\tau} \right) dY \tag{1}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

• If $|h^{1/2}Y_1| \ge \epsilon$

Recall we are estimating

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L^{2}_{t,x}}^{2} \lesssim h^{1-2s}\int_{-T/h}^{T/h}\int\langle h^{\frac{1}{2}}Y_{1}\langle\tau\rangle+z_{1}+\tau\rangle^{-2\nu}\exp\left(-Y^{2}\right)dYd\tau.$$

▶ In the region $|h^{1/2}Y_1| \le \epsilon$ ($\epsilon \ll 1$ fixed), we integrate in time by using the variable

$$ilde{ au} = au + h^{rac{1}{2}} Y_1 \langle au
angle \qquad ({ t Jacobian} = 1 + O(\epsilon))$$

so we bound the integral by

$$h^{1-2s} \int \left(\int_{-CT/h}^{CT/h} \langle z_1 + \tilde{\tau} \rangle^{-2\nu} \exp\left(-Y^2\right) d\tilde{\tau} \right) dY \tag{1}$$

► If $|h^{1/2}Y_1| \ge \epsilon$, then $|Y_1| \gtrsim h^{-\frac{1}{2}}$ and $\langle h^{\frac{1}{2}}Y_1\langle \tau \rangle + z_1 + \tau \rangle^{-2\nu} \exp\left(-Y^2\right) \lesssim \langle z_1 + \tau \rangle^{-2\nu} \exp\left(-Y^2/2\right) O(h^{\infty})$ ⇒ Integral $\le (1) \times O(h^{\infty})$

Recall we are estimating

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L^{2}_{t,x}}^{2} \lesssim h^{1-2s}\int_{-T/h}^{T/h}\int\langle h^{\frac{1}{2}}Y_{1}\langle\tau\rangle+z_{1}+\tau\rangle^{-2\nu}\exp\left(-Y^{2}\right)dYd\tau.$$

▶ In the region $|h^{1/2}Y_1| \le \epsilon$ ($\epsilon \ll 1$ fixed), we integrate in time by using the variable

$$ilde{ au} = au + h^{rac{1}{2}} \, Y_1 \langle au
angle \qquad ({ t Jacobian} = 1 + O(\epsilon))$$

so we bound the integral by

$$h^{1-2s} \int \left(\int_{-CT/h}^{CT/h} \langle z_1 + \tilde{\tau} \rangle^{-2\nu} \exp\left(-Y^2\right) d\tilde{\tau} \right) dY \tag{1}$$

• If $|h^{1/2}Y_1| \ge \epsilon$, then $|Y_1| \gtrsim h^{-rac{1}{2}}$ and

$$\langle h^{rac{1}{2}} Y_1 \langle au
angle + z_1 + au
angle^{-2
u} \exp\left(-Y^2
ight) \lesssim \langle z_1 + au
angle^{-2
u} \exp\left(-Y^2/2
ight) O(h^\infty)$$

 \Rightarrow Integral \leq (1) imes $\mathit{O}(h^{\infty})$

Conclusion: If $s = \frac{1}{2}$ and $\nu > \frac{1}{2}$

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{\frac{1}{2}}G^{t}_{z,\zeta,h}\right|\right|_{L^{2}\left([-T,T]\times\mathbb{R}^{n}\right)}\leq C$$

Recall we are estimating

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L^{2}_{t,x}}^{2} \lesssim h^{1-2s}\int_{-T/h}^{T/h}\int\langle h^{\frac{1}{2}}Y_{1}\langle\tau\rangle+z_{1}+\tau\rangle^{-2\nu}\exp\left(-Y^{2}\right)dYd\tau.$$

▶ In the region $|h^{1/2}Y_1| \le \epsilon$ ($\epsilon \ll 1$ fixed), we integrate in time by using the variable

$$ilde{ au} = au + h^{rac{1}{2}} Y_1 \langle au
angle \qquad ({ t Jacobian} = 1 + O(\epsilon))$$

so we bound the integral by

$$h^{1-2s} \int \left(\int_{-CT/h}^{CT/h} \langle z_1 + \tilde{\tau} \rangle^{-2\nu} \exp\left(-Y^2\right) d\tilde{\tau} \right) dY \tag{1}$$

• If $|h^{1/2}Y_1| \ge \epsilon$, then $|Y_1| \gtrsim h^{-\frac{1}{2}}$ and $\langle h^{\frac{1}{2}}Y_1\langle \tau \rangle + z_1 + \tau \rangle^{-2\nu} \exp(-Y^2) \lesssim \langle z_1 + \tau \rangle^{-2\nu} \exp(-Y^2/2) O(h^{\infty})$ \Rightarrow Integral $\le (1) \times O(h^{\infty})$ Conclusion: If $s = \frac{1}{2}$ and $\nu > \frac{1}{2}$

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{\frac{1}{2}}G^{t}_{z,\zeta,h}\right|\right|_{L^{2}\left([-T,T]\times\mathbb{R}^{n}\right)}\leq C \qquad \text{uniformly in } h\in(0,1]$$

Recall we are estimating

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{s}G_{z,\zeta,h}^{t}\right|\right|_{L^{2}_{t,x}}^{2} \lesssim h^{1-2s}\int_{-T/h}^{T/h}\int\langle h^{\frac{1}{2}}Y_{1}\langle\tau\rangle+z_{1}+\tau\rangle^{-2\nu}\exp\left(-Y^{2}\right)dYd\tau.$$

▶ In the region $|h^{1/2}Y_1| \le \epsilon$ ($\epsilon \ll 1$ fixed), we integrate in time by using the variable

$$ilde{ au} = au + h^{rac{1}{2}} Y_1 \langle au
angle \qquad ({ t Jacobian} = 1 + O(\epsilon))$$

so we bound the integral by

$$h^{1-2s} \int \left(\int_{-CT/h}^{CT/h} \langle z_1 + \tilde{\tau} \rangle^{-2\nu} \exp\left(-Y^2\right) d\tilde{\tau} \right) dY \tag{1}$$

▶ If $|h^{1/2}Y_1| \ge \epsilon$, then $|Y_1| \gtrsim h^{-\frac{1}{2}}$ and $\langle h^{\frac{1}{2}}Y_1\langle \tau \rangle + z_1 + \tau \rangle^{-2\nu} \exp\left(-Y^2\right) \lesssim \langle z_1 + \tau \rangle^{-2\nu} \exp\left(-Y^2/2\right) O(h^{\infty})$

 $\Rightarrow \text{Integral} \leq (1) \times O(h^{\infty})$ Conclusion: If $s = \frac{1}{2}$ and $\nu > \frac{1}{2}$

$$\left|\left|\langle x\rangle^{-\nu}\langle\zeta/h\rangle^{\frac{1}{2}}G^{t}_{z,\zeta,h}\right|\right|_{L^{2}\left([-\tau,\tau]\times\mathbb{R}^{n}\right)}\leq C$$

uniformly in $h \in (0, 1]$ and in $z \in \mathbb{R}^n$.

General problem: Extend Strichartz estimates to asymptotically flat manifolds

General problem: Extend Strichartz estimates to asymptotically flat manifolds

 $1. \ \mbox{see}$ which properties persist or can be lost

General problem: Extend Strichartz estimates to asymptotically flat manifolds

(ロ)、(型)、(E)、(E)、 E) の(の)

- 1. see which properties persist or can be lost
- 2. more specifically, try to decouple what happens near infinity

General problem: Extend Strichartz estimates to asymptotically flat manifolds

- $1. \ {\rm see}$ which properties persist or can be lost
- 2. more specifically, try to decouple what happens near infinity (where one expects the same behavior as on $\mathbb{R}^n)$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ □臣 = のへで

General problem: Extend Strichartz estimates to asymptotically flat manifolds

- 1. see which properties persist or can be lost
- 2. more specifically, try to decouple what happens near infinity (where one expects the same behavior as on \mathbb{R}^n) from what happens inside a compact set

General problem: Extend Strichartz estimates to asymptotically flat manifolds

- 1. see which properties persist or can be lost
- 2. more specifically, try to decouple what happens near infinity (where one expects the same behavior as on \mathbb{R}^n) from what happens inside a compact set (where the geometry/geodesic flow may be arbitrary and complicated)

General problem: Extend Strichartz estimates to asymptotically flat manifolds

- 1. see which properties persist or can be lost
- 2. more specifically, try to decouple what happens near infinity (where one expects the same behavior as on \mathbb{R}^n) from what happens inside a compact set (where the geometry/geodesic flow may be arbitrary and complicated)

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● のへで

3. see the influence of the geometry on nonlinear equations

General problem: Extend Strichartz estimates to asymptotically flat manifolds

- 1. see which properties persist or can be lost
- 2. more specifically, try to decouple what happens near infinity (where one expects the same behavior as on \mathbb{R}^n) from what happens inside a compact set (where the geometry/geodesic flow may be arbitrary and complicated)

- 3. see the influence of the geometry on nonlinear equations
- 4. the Schrödinger equation can be replaced by other dispersive PDE (wave, Klein-Gordon) which are relevant on asymptotically flat manifolds

General problem: Extend Strichartz estimates to asymptotically flat manifolds

- 1. see which properties persist or can be lost
- more specifically, try to decouple what happens near infinity (where one expects the same behavior as on Rⁿ) from what happens inside a compact set (where the geometry/geodesic flow may be arbitrary and complicated)
- 3. see the influence of the geometry on nonlinear equations
- 4. the Schrödinger equation can be replaced by other dispersive PDE (wave, Klein-Gordon) which are relevant on asymptotically flat manifolds
- 5. good motivation / test to understand which scattering properties are robust and relevant (in particular in the low energy analysis)

General problem: Extend Strichartz estimates to asymptotically flat manifolds

- 1. see which properties persist or can be lost
- more specifically, try to decouple what happens near infinity (where one expects the same behavior as on Rⁿ) from what happens inside a compact set (where the geometry/geodesic flow may be arbitrary and complicated)
- 3. see the influence of the geometry on nonlinear equations
- 4. the Schrödinger equation can be replaced by other dispersive PDE (wave, Klein-Gordon) which are relevant on asymptotically flat manifolds
- 5. good motivation / test to understand which scattering properties are robust and relevant (in particular in the low energy analysis)

Scattering inequalities turn out to play a crucial role in this problem.

• The model: \mathbb{R}^n , equipped with the flat metric,

$$G_0 = dx_1^2 + \cdots + dx_n^2 = \sum_{j,k} G_{jk} dx_j dx_k, \qquad G_0 := (G_{jk}) = I.$$

▶ The model: \mathbb{R}^n , equipped with the flat metric,

$$G_0 = dx_1^2 + \cdots + dx_n^2 = \sum_{j,k} G_{jk} dx_j dx_k, \qquad G_0 := (G_{jk}) = I.$$

The geodesic flow $\phi^t : \mathbb{R}^n \times \mathbb{R}^n$

▶ The model: \mathbb{R}^n , equipped with the flat metric,

$$G_0 = dx_1^2 + \cdots + dx_n^2 = \sum_{j,k} G_{jk} dx_j dx_k, \qquad G_0 := (G_{jk}) = I.$$

The geodesic flow $\phi^t : \mathbb{R}^n \times \mathbb{R}^n (= T^* \mathbb{R}^n)$

▶ The model: \mathbb{R}^n , equipped with the flat metric,

$$G_0 = dx_1^2 + \cdots + dx_n^2 = \sum_{j,k} G_{jk} dx_j dx_k, \qquad G_0 := (G_{jk}) = I.$$

The geodesic flow $\phi^t : \mathbb{R}^n \times \mathbb{R}^n (= T^* \mathbb{R}^n) \to \mathbb{R}^n \times \mathbb{R}^n$ is given by

$$\phi^t(x,\xi) = (x+2t\xi,\xi)$$

• The model: \mathbb{R}^n , equipped with the flat metric,

$$G_0 = dx_1^2 + \cdots + dx_n^2 = \sum_{j,k} G_{jk} dx_j dx_k, \qquad G_0 := (G_{jk}) = I.$$

The geodesic flow $\phi^t : \mathbb{R}^n \times \mathbb{R}^n (= T^* \mathbb{R}^n) \to \mathbb{R}^n \times \mathbb{R}^n$ is given by

$$\phi^t(x,\xi) = (x+2t\xi,\xi) =: (x^t,\xi^t),$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

▶ The model: \mathbb{R}^n , equipped with the flat metric,

$$G_0 = dx_1^2 + \cdots + dx_n^2 = \sum_{j,k} G_{jk} dx_j dx_k, \qquad G_0 := (G_{jk}) = I.$$

The geodesic flow $\phi^t : \mathbb{R}^n \times \mathbb{R}^n (= T^* \mathbb{R}^n) \to \mathbb{R}^n \times \mathbb{R}^n$ is given by

$$\phi^t(x,\xi) = (x+2t\xi,\xi) =: (x^t,\xi^t),$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

it solves the Hamilton equations

▶ The model: \mathbb{R}^n , equipped with the flat metric,

$$G_0 = dx_1^2 + \cdots + dx_n^2 = \sum_{j,k} G_{jk} dx_j dx_k, \qquad G_0 := (G_{jk}) = I.$$

The geodesic flow $\phi^t: \mathbb{R}^n \times \mathbb{R}^n (= T^* \mathbb{R}^n) \to \mathbb{R}^n \times \mathbb{R}^n$ is given by

$$\phi^t(x,\xi) = (x+2t\xi,\xi) =: (x^t,\xi^t),$$

it solves the Hamilton equations

$$\dot{x}^t = (\partial_{\xi} p)(x^t, \xi^t), \qquad \dot{\xi}^t = -(\partial_x p)(x^t, \xi^t)$$

where

$$p(x,\xi) = |\xi|^2 = \xi \cdot G_0^{-1}\xi$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

▶ The model: \mathbb{R}^n , equipped with the flat metric,

$$G_0 = dx_1^2 + \dots + dx_n^2 = \sum_{j,k} G_{jk} dx_j dx_k, \qquad G_0 := (G_{jk}) = I.$$

The geodesic flow $\phi^t: \mathbb{R}^n \times \mathbb{R}^n (= T^* \mathbb{R}^n) \to \mathbb{R}^n \times \mathbb{R}^n$ is given by

$$\phi^t(x,\xi) = (x+2t\xi,\xi) =: (x^t,\xi^t),$$

it solves the Hamilton equations

$$\dot{x}^t = (\partial_{\xi} p)(x^t, \xi^t), \qquad \dot{\xi}^t = -(\partial_x p)(x^t, \xi^t)$$

where

$$p(x,\xi) = |\xi|^2 = \xi \cdot G_0^{-1}\xi$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

is the (principal) symbol of $-\Delta = D_1^2 + \cdots + D_n^2$ with $D_j = \frac{1}{i} \frac{\partial}{\partial x_j}$

▶ The model: \mathbb{R}^n , equipped with the flat metric,

$$G_0 = dx_1^2 + \cdots + dx_n^2 = \sum_{j,k} G_{jk} dx_j dx_k, \qquad G_0 := (G_{jk}) = I.$$

The geodesic flow $\phi^t : \mathbb{R}^n \times \mathbb{R}^n (= T^* \mathbb{R}^n) \to \mathbb{R}^n \times \mathbb{R}^n$ is given by

$$\phi^t(x,\xi) = (x+2t\xi,\xi) =: (x^t,\xi^t),$$

it solves the Hamilton equations

$$\dot{x}^t = (\partial_{\xi} p)(x^t, \xi^t), \qquad \dot{\xi}^t = -(\partial_x p)(x^t, \xi^t)$$

where

$$p(x,\xi) = |\xi|^2 = \xi \cdot G_0^{-1}\xi$$

is the (principal) symbol of $-\Delta = D_1^2 + \cdots + D_n^2$ with $D_j = \frac{1}{i} \frac{\partial}{\partial x_j}$

▶ Pertubed model: \mathbb{R}^n , equipped with a metric $\sum_{i,k} G_{jk}(x) dx_j dx_k$ such that

$$G(x) - I
ightarrow 0$$
 as $x
ightarrow \infty$, $G(x) := (G_{jk}(x))$

▶ The model: \mathbb{R}^n , equipped with the flat metric,

$$G_0 = dx_1^2 + \cdots + dx_n^2 = \sum_{j,k} G_{jk} dx_j dx_k, \qquad G_0 := (G_{jk}) = I.$$

The geodesic flow $\phi^t : \mathbb{R}^n \times \mathbb{R}^n (= T^* \mathbb{R}^n) \to \mathbb{R}^n \times \mathbb{R}^n$ is given by

$$\phi^t(x,\xi) = (x+2t\xi,\xi) =: (x^t,\xi^t),$$

it solves the Hamilton equations

$$\dot{x}^t = (\partial_{\xi} p)(x^t, \xi^t), \qquad \dot{\xi}^t = -(\partial_x p)(x^t, \xi^t)$$

where

$$p(x,\xi) = |\xi|^2 = \xi \cdot G_0^{-1}\xi$$

is the (principal) symbol of $-\Delta = D_1^2 + \cdots + D_n^2$ with $D_j = \frac{1}{i} \frac{\partial}{\partial x_j}$

▶ Pertubed model: \mathbb{R}^n , equipped with a metric $\sum_{i,k} G_{jk}(x) dx_j dx_k$ such that

$$G(x) - I
ightarrow 0$$
 as $x
ightarrow \infty$, $G(x) := (G_{jk}(x))$

more precisely, $\partial^{\alpha}(G_{jk}(x) - \delta_{jk}) = O(\langle x \rangle^{-\mu - |\alpha|})$ for some $\mu > 0$.

▶ The model: \mathbb{R}^n , equipped with the flat metric,

$$G_0 = dx_1^2 + \cdots + dx_n^2 = \sum_{j,k} G_{jk} dx_j dx_k, \qquad G_0 := (G_{jk}) = I.$$

The geodesic flow $\phi^t : \mathbb{R}^n \times \mathbb{R}^n (= T^* \mathbb{R}^n) \to \mathbb{R}^n \times \mathbb{R}^n$ is given by

$$\phi^t(x,\xi) = (x+2t\xi,\xi) =: (x^t,\xi^t),$$

it solves the Hamilton equations

$$\dot{x}^t = (\partial_{\xi} p)(x^t, \xi^t), \qquad \dot{\xi}^t = -(\partial_x p)(x^t, \xi^t)$$

where

$$p(x,\xi) = |\xi|^2 = \xi \cdot G_0^{-1}\xi$$

is the (principal) symbol of $-\Delta = D_1^2 + \cdots + D_n^2$ with $D_j = \frac{1}{i} \frac{\partial}{\partial x_j}$

▶ Pertubed model: \mathbb{R}^n , equipped with a metric $\sum_{j,k} G_{jk}(x) dx_j dx_k$ such that

$$G(x) - I
ightarrow 0$$
 as $x
ightarrow \infty$, $G(x) := (G_{jk}(x))$

more precisely, $\partial^{\alpha}(G_{jk}(x) - \delta_{jk}) = O(\langle x \rangle^{-\mu - |\alpha|})$ for some $\mu > 0$. The geodesic flow is defined analogously with

$$p(x,\xi) = \xi \cdot G(x)^{-1}\xi$$

▶ The model: \mathbb{R}^n , equipped with the flat metric,

$$G_0 = dx_1^2 + \cdots + dx_n^2 = \sum_{j,k} G_{jk} dx_j dx_k, \qquad G_0 := (G_{jk}) = I.$$

The geodesic flow $\phi^t : \mathbb{R}^n \times \mathbb{R}^n (= T^* \mathbb{R}^n) \to \mathbb{R}^n \times \mathbb{R}^n$ is given by

$$\phi^t(x,\xi) = (x+2t\xi,\xi) =: (x^t,\xi^t),$$

it solves the Hamilton equations

$$\dot{x}^t = (\partial_{\xi} p)(x^t, \xi^t), \qquad \dot{\xi}^t = -(\partial_x p)(x^t, \xi^t)$$

where

$$p(x,\xi) = |\xi|^2 = \xi \cdot G_0^{-1}\xi$$

is the (principal) symbol of $-\Delta = D_1^2 + \cdots + D_n^2$ with $D_j = \frac{1}{i} \frac{\partial}{\partial x_j}$

▶ Pertubed model: \mathbb{R}^n , equipped with a metric $\sum_{i,k} G_{ik}(x) dx_i dx_k$ such that

$$G(x) - I \rightarrow 0$$
 as $x \rightarrow \infty$, $G(x) := (G_{jk}(x))$

more precisely, $\partial^{\alpha}(G_{jk}(x) - \delta_{jk}) = O(\langle x \rangle^{-\mu - |\alpha|})$ for some $\mu > 0$. The geodesic flow is defined analogously with

$$p(x,\xi) = \xi \cdot G(x)^{-1}\xi = \sum_{j,k} G^{jk}(x)\xi_j\xi_k$$

▶ The model: \mathbb{R}^n , equipped with the flat metric,

$$G_0 = dx_1^2 + \cdots + dx_n^2 = \sum_{j,k} G_{jk} dx_j dx_k, \qquad G_0 := (G_{jk}) = I.$$

The geodesic flow $\phi^t : \mathbb{R}^n \times \mathbb{R}^n (= T^* \mathbb{R}^n) \to \mathbb{R}^n \times \mathbb{R}^n$ is given by

$$\phi^t(x,\xi) = (x+2t\xi,\xi) =: (x^t,\xi^t),$$

it solves the Hamilton equations

$$\dot{x}^t = (\partial_{\xi} p)(x^t, \xi^t), \qquad \dot{\xi}^t = -(\partial_x p)(x^t, \xi^t)$$

where

$$p(x,\xi) = |\xi|^2 = \xi \cdot G_0^{-1}\xi$$

is the (principal) symbol of $-\Delta = D_1^2 + \cdots + D_n^2$ with $D_j = \frac{1}{i} \frac{\partial}{\partial x_j}$

▶ Pertubed model: \mathbb{R}^n , equipped with a metric $\sum_{i,k} G_{ik}(x) dx_i dx_k$ such that

$$G(x) - I \rightarrow 0$$
 as $x \rightarrow \infty$, $G(x) := (G_{jk}(x))$

more precisely, $\partial^{\alpha}(G_{jk}(x) - \delta_{jk}) = O(\langle x \rangle^{-\mu - |\alpha|})$ for some $\mu > 0$. The geodesic flow is defined analogously with

$$p(x,\xi) = \xi \cdot G(x)^{-1}\xi = \sum_{j,k} G^{jk}(x)\xi_j\xi_k$$

the (principal) symbol of the Laplace-Beltrami operator

$$-\Delta_{G} = -\sum_{j,k} G^{jk}(x) \partial_{x_{j}} \partial_{x_{k}} + \sum_{j,k,\ell} G^{jk}(x) \Gamma_{jk}^{\ell}(x) \partial_{x_{\ell}}$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● のへで

More general model: asymptotically conical manifolds.

◆□ ▶ < 圖 ▶ < 圖 ▶ < 圖 ▶ < 圖 • 의 Q @</p>

• More general model: asymptotically conical manifolds. In polar coordinates, $\mathbb{R}^n \setminus 0$ equipped with the Euclidean metric is isometric to

 $(0,+\infty) imes\mathbb{S}^{n-1}$ equipped with $dr^2+r^2g_{\mathbb{S}^{n-1}}$

with $g_{\mathbb{S}^{n-1}}$ the standard metric on the sphere.

• More general model: asymptotically conical manifolds. In polar coordinates, $\mathbb{R}^n \setminus 0$ equipped with the Euclidean metric is isometric to

$$(0,+\infty) imes \mathbb{S}^{n-1}$$
 equipped with $dr^2+r^2g_{\mathbb{S}^{n-1}}$

with $g_{\mathbb{S}^{n-1}}$ the standard metric on the sphere.

An asymptotically conical manifold is of the form $M=M_{\rm c}\sqcup M_{\infty}$ with $M_{\rm c}$ compact with boundary

 $M_{\infty} \approx (R,\infty)_r \times S$ equipped with $G = dr^2 + r^2 g(r)$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

with S compact (without boundary), $\dim(S) = n - 1$,

More general model: asymptotically conical manifolds.
 In polar coordinates, Rⁿ \ 0 equipped with the Euclidean metric is isometric to

$$(0,+\infty) imes \mathbb{S}^{n-1}$$
 equipped with $dr^2+r^2g_{\mathbb{S}^{n-1}}$

with $g_{\mathbb{S}^{n-1}}$ the standard metric on the sphere. An asymptotically conical manifold is of the form $M = M_c \sqcup M_\infty$ with M_c compact with boundary

 $M_{\infty} \approx (R, \infty)_r \times S$ equipped with $G = dr^2 + r^2 g(r)$

with S compact (without boundary), dim(S) = n-1, and, for some metric $g(\infty)$ on S and some $\mu \in (0,1]$,

$$\partial_r^k (g(r) - g(\infty)) = O(\langle r \rangle^{-\mu-k}).$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

More general model: asymptotically conical manifolds.
 In polar coordinates, Rⁿ \ 0 equipped with the Euclidean metric is isometric to

$$(0,+\infty) imes \mathbb{S}^{n-1}$$
 equipped with $dr^2+r^2g_{\mathbb{S}^{n-1}}$

with $g_{\mathbb{S}^{n-1}}$ the standard metric on the sphere. An asymptotically conical manifold is of the form $M = M_c \sqcup M_\infty$ with M_c compact with boundary

 $M_{\infty} \approx (R, \infty)_r \times S$ equipped with $G = dr^2 + r^2 g(r)$

with S compact (without boundary), dim(S) = n-1, and, for some metric $g(\infty)$ on S and some $\mu \in (0,1]$,

$$\partial_r^k (g(r) - g(\infty)) = O(\langle r \rangle^{-\mu-k}).$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● ● ●

More general model: asymptotically conical manifolds.
 In polar coordinates, Rⁿ \ 0 equipped with the Euclidean metric is isometric to

$$(0,+\infty) imes \mathbb{S}^{n-1}$$
 equipped with $dr^2+r^2g_{\mathbb{S}^{n-1}}$

with $g_{\mathbb{S}^{n-1}}$ the standard metric on the sphere. An asymptotically conical manifold is of the form $M = M_c \sqcup M_\infty$ with M_c compact with boundary

 $M_{\infty} \approx (R, \infty)_r \times S$ equipped with $G = dr^2 + r^2 g(r)$

with S compact (without boundary), dim(S) = n-1, and, for some metric $g(\infty)$ on S and some $\mu \in (0,1]$,

$$\partial_r^k (g(r) - g(\infty)) = O(\langle r \rangle^{-\mu-k}).$$

Motivation to study such models:

More general model: asymptotically conical manifolds.
 In polar coordinates, Rⁿ \ 0 equipped with the Euclidean metric is isometric to

$$(0,+\infty) imes\mathbb{S}^{n-1}$$
 equipped with $dr^2+r^2g_{\mathbb{S}^{n-1}}$

with $g_{\mathbb{S}^{n-1}}$ the standard metric on the sphere. An asymptotically conical manifold is of the form $M = M_c \sqcup M_\infty$ with M_c compact with boundary

 $M_{\infty} \approx (R, \infty)_r \times S$ equipped with $G = dr^2 + r^2 g(r)$

with S compact (without boundary), dim(S) = n-1, and, for some metric $g(\infty)$ on S and some $\mu \in (0,1]$,

$$\partial_r^k (g(r) - g(\infty)) = O(\langle r \rangle^{-\mu-k}).$$

Motivation to study such models:

Good models of scattering theory

More general model: asymptotically conical manifolds.
 In polar coordinates, Rⁿ \ 0 equipped with the Euclidean metric is isometric to

$$(0,+\infty) imes\mathbb{S}^{n-1}$$
 equipped with $dr^2+r^2g_{\mathbb{S}^{n-1}}$

with $g_{\mathbb{S}^{n-1}}$ the standard metric on the sphere. An asymptotically conical manifold is of the form $M = M_c \sqcup M_\infty$ with M_c compact with boundary

 $M_{\infty} \approx (R, \infty)_r \times S$ equipped with $G = dr^2 + r^2 g(r)$

with S compact (without boundary), dim(S) = n-1, and, for some metric $g(\infty)$ on S and some $\mu \in (0,1]$,

$$\partial_r^k (g(r) - g(\infty)) = O(\langle r \rangle^{-\mu-k}).$$

Motivation to study such models:

- Good models of scattering theory
- time slices of certain space-times
Asymptotically flat manifolds

More general model: asymptotically conical manifolds.
 In polar coordinates, Rⁿ \ 0 equipped with the Euclidean metric is isometric to

$$(0,+\infty) imes\mathbb{S}^{n-1}$$
 equipped with $dr^2+r^2g_{\mathbb{S}^{n-1}}$

with $g_{\mathbb{S}^{n-1}}$ the standard metric on the sphere. An asymptotically conical manifold is of the form $M = M_c \sqcup M_\infty$ with M_c compact with boundary

$$M_{\infty} \approx (R,\infty)_r \times S$$
 equipped with $G = dr^2 + r^2 g(r)$

with S compact (without boundary), dim(S) = n-1, and, for some metric $g(\infty)$ on S and some $\mu \in (0,1]$,

$$\partial_r^k (g(r) - g(\infty)) = O(\langle r \rangle^{-\mu-k}).$$

Motivation to study such models:

- Good models of scattering theory
- time slices of certain space-times
- allow to describe the propagation into an inhomogeneous medium, with possible impurities (small perturbations) at infinity and strong perturbation inside a compact set

Asymptotically flat manifolds

More general model: asymptotically conical manifolds.
 In polar coordinates, Rⁿ \ 0 equipped with the Euclidean metric is isometric to

$$(0,+\infty) imes\mathbb{S}^{n-1}$$
 equipped with $dr^2+r^2g_{\mathbb{S}^{n-1}}$

with $g_{\mathbb{S}^{n-1}}$ the standard metric on the sphere. An asymptotically conical manifold is of the form $M = M_c \sqcup M_\infty$ with M_c compact with boundary

$$M_{\infty} \approx (R,\infty)_r \times S$$
 equipped with $G = dr^2 + r^2 g(r)$

with S compact (without boundary), dim(S) = n-1, and, for some metric $g(\infty)$ on S and some $\mu \in (0,1]$,

$$\partial_r^k (g(r) - g(\infty)) = O(\langle r \rangle^{-\mu-k}).$$

Motivation to study such models:

- Good models of scattering theory
- time slices of certain space-times
- allow to describe the propagation into an inhomogeneous medium, with possible impurities (small perturbations) at infinity and strong perturbation inside a compact set

Let P be the selfadjoint realization of $-\Delta_G$ on $L^2(M)$, with (M, G) an asymptotically flat manifold. We let

$$R(z) = (P - z)^{-1}, \qquad z \in \mathbb{C} \setminus [0, +\infty)$$

・ロト・日本・モト・モート ヨー うへで

Let P be the selfadjoint realization of $-\Delta_G$ on $L^2(M)$, with (M, G) an asymptotically flat manifold. We let

$$R(z) = (P-z)^{-1}, \qquad z \in \mathbb{C} \setminus [0, +\infty)$$

・ロト・日本・モト・モート ヨー うへで

Rem: recall that spec(P) $\subset [0, \infty)$ since $(Pu, u)_{L^2} = |||\nabla_G u|||_{L^2}^2 \ge 0$

Let P be the selfadjoint realization of $-\Delta_G$ on $L^2(M)$, with (M, G) an asymptotically flat manifold. We let

$$R(z) = (P-z)^{-1}, \qquad z \in \mathbb{C} \setminus [0, +\infty)$$

Rem: recall that spec(P) \subset [0, ∞) since (Pu, u)_{L²} = $|||\nabla_G u|||_{L²}^2 \ge 0$ **Facts**:

 P has no (embbeded) eigenvalues, i.e. the spectrum is continuous (Froese-Herbst 82, Donnelly 99, Koch-Tataru 06, Ito-Skibsted 13)

Let P be the selfadjoint realization of $-\Delta_G$ on $L^2(M)$, with (M, G) an asymptotically flat manifold. We let

$$R(z) = (P-z)^{-1}, \qquad z \in \mathbb{C} \setminus [0, +\infty)$$

Rem: recall that spec(P) \subset [0, ∞) since (Pu, u)_{L²} = $|||\nabla_G u|||_{L²}^2 \ge 0$ **Facts**:

- P has no (embbeded) eigenvalues, i.e. the spectrum is continuous (Froese-Herbst 82, Donnelly 99, Koch-Tataru 06, Ito-Skibsted 13)
- there is a limiting absorption principle, i.e.

$$\langle r \rangle^{-\nu} R(\lambda \pm i0) \langle r \rangle^{-\nu} : L^2(M) \to L^2(M)$$

(日) (日) (日) (日) (日) (日) (日) (日)

exists at all positive energies if $\nu > \frac{1}{2}$,

Let P be the selfadjoint realization of $-\Delta_G$ on $L^2(M)$, with (M, G) an asymptotically flat manifold. We let

$$R(z) = (P-z)^{-1}, \qquad z \in \mathbb{C} \setminus [0, +\infty)$$

Rem: recall that spec(P) \subset [0, ∞) since (Pu, u)_{L²} = $|||\nabla_G u|||_{L²}^2 \ge 0$ **Facts**:

- P has no (embbeded) eigenvalues, i.e. the spectrum is continuous (Froese-Herbst 82, Donnelly 99, Koch-Tataru 06, Ito-Skibsted 13)
- there is a limiting absorption principle, i.e.

$$\langle r \rangle^{-\nu} R(\lambda \pm i0) \langle r \rangle^{-\nu} : L^2(M) \to L^2(M)$$

(日) (日) (日) (日) (日) (日) (日) (日)

exists at all positive energies if $\nu > \frac{1}{2}$, and is C^k on $(0, \infty)$ if $\nu > \frac{1}{2} + k$

Let P be the selfadjoint realization of $-\Delta_G$ on $L^2(M)$, with (M, G) an asymptotically flat manifold. We let

$$R(z) = (P-z)^{-1}, \qquad z \in \mathbb{C} \setminus [0, +\infty)$$

Rem: recall that spec(P) \subset [0, ∞) since (Pu, u)_{L²} = $|||\nabla_G u|||_{L²}^2 \ge 0$ **Facts**:

- P has no (embbeded) eigenvalues, i.e. the spectrum is continuous (Froese-Herbst 82, Donnelly 99, Koch-Tataru 06, Ito-Skibsted 13)
- there is a limiting absorption principle, i.e.

$$\langle r \rangle^{-\nu} R(\lambda \pm i0) \langle r \rangle^{-\nu} : L^2(M) \to L^2(M)$$

exists at all positive energies if $\nu > \frac{1}{2}$, and is C^k on $(0, \infty)$ if $\nu > \frac{1}{2} + k$ (consequence of the Mourre Theory, [Jensen-Mourre-Perry])

Let P be the selfadjoint realization of $-\Delta_G$ on $L^2(M)$, with (M, G) an asymptotically flat manifold. We let

$$R(z) = (P-z)^{-1}, \qquad z \in \mathbb{C} \setminus [0, +\infty)$$

Rem: recall that spec(P) \subset [0, ∞) since (Pu, u)_{L²} = $|||\nabla_G u|||_{L²}^2 \ge 0$ **Facts**:

- P has no (embbeded) eigenvalues, i.e. the spectrum is continuous (Froese-Herbst 82, Donnelly 99, Koch-Tataru 06, Ito-Skibsted 13)
- there is a limiting absorption principle, i.e.

$$\langle r \rangle^{-\nu} R(\lambda \pm i0) \langle r \rangle^{-\nu} : L^2(M) \to L^2(M)$$

exists at all positive energies if $\nu > \frac{1}{2}$, and is C^k on $(0, \infty)$ if $\nu > \frac{1}{2} + k$ (consequence of the Mourre Theory, [Jensen-Mourre-Perry])

▶ In particular, for any $\varphi \in C_0^\infty(0, +\infty)$ and $\lambda > 0$,

$$\left|\left|\langle r\rangle^{-\nu}\varphi(P/\lambda)e^{-itP}\langle r\rangle^{-\nu}\right|\right|_{L^{2}(M)\to L^{2}(M)} \leq C_{\lambda,\varphi,\nu}\langle t\rangle^{-k}$$
(2)

if $\nu > \frac{1}{2} + k$

Let P be the selfadjoint realization of $-\Delta_G$ on $L^2(M)$, with (M, G) an asymptotically flat manifold. We let

$$R(z) = (P - z)^{-1}, \qquad z \in \mathbb{C} \setminus [0, +\infty)$$

Rem: recall that spec(P) \subset [0, ∞) since (Pu, u)_{L²} = $|||\nabla_G u|||_{L²}^2 \ge 0$ **Facts**:

- P has no (embbeded) eigenvalues, i.e. the spectrum is continuous (Froese-Herbst 82, Donnelly 99, Koch-Tataru 06, Ito-Skibsted 13)
- there is a limiting absorption principle, i.e.

$$\langle r \rangle^{-\nu} R(\lambda \pm i0) \langle r \rangle^{-\nu} : L^2(M) \to L^2(M)$$

exists at all positive energies if $\nu > \frac{1}{2}$, and is C^k on $(0, \infty)$ if $\nu > \frac{1}{2} + k$ (consequence of the Mourre Theory, [Jensen-Mourre-Perry])

▶ In particular, for any $\varphi \in C_0^\infty(0, +\infty)$ and $\lambda > 0$,

$$\left|\left|\langle r\rangle^{-\nu}\varphi(P/\lambda)e^{-itP}\langle r\rangle^{-\nu}\right|\right|_{L^{2}(M)\to L^{2}(M)} \leq C_{\lambda,\varphi,\nu}\langle t\rangle^{-k}$$
(2)

if $\nu > \frac{1}{2} + k$

Question: behavior of $R(\lambda \pm i0)$ and (2) as $\lambda \to \infty$ (high energy) and $\lambda \to 0$ (low energy) ?

Scattering estimates on asymptotically flat manifolds High energy estimates $(\lambda \rightarrow +\infty)$

◆□ → <圖 → < Ξ → < Ξ → < Ξ · 9 < @</p>

High energy estimates $(\lambda \to +\infty)$ depend on the behavior of the geodesic flow ϕ^t

・ロト・日本・モト・モート ヨー うへで

High energy estimates $(\lambda
ightarrow +\infty)$ depend on the behavior of the geodesic flow ϕ^t

Worst case: general case

High energy estimates $(\lambda
ightarrow +\infty)$ depend on the behavior of the geodesic flow ϕ^t

• Worst case: general case (everywhere below $\nu > 1/2$)

$$\left|\left|\langle r\rangle^{-\nu}R(\lambda\pm i0)\langle r\rangle^{-\nu}\right|\right|_{L^{2}(M)\to L^{2}(M)}\lesssim e^{C\lambda^{1/2}}$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

[Burq , Cardoso-Vodev]

High energy estimates $(\lambda \to +\infty)$ depend on the behavior of the geodesic flow ϕ^t

• Worst case: general case (everywhere below $\nu > 1/2$)

$$\left|\left|\langle r\rangle^{-\nu}R(\lambda\pm i0)\langle r\rangle^{-\nu}\right|\right|_{L^{2}(M)\to L^{2}(M)}\lesssim e^{C\lambda^{1/2}}$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

[Burq , Cardoso-Vodev]

Best case: non trapping geodesic flow

High energy estimates $(\lambda \to +\infty)$ depend on the behavior of the geodesic flow ϕ^t

• Worst case: general case (everywhere below $\nu > 1/2$)

$$\left|\left|\langle r\rangle^{-\nu}R(\lambda\pm i0)\langle r\rangle^{-\nu}\right|\right|_{L^{2}(M)\to L^{2}(M)}\lesssim e^{C\lambda^{1/2}}$$

[Burq , Cardoso-Vodev]

Best case: non trapping geodesic flow

$$\left|\left|\langle r\rangle^{-\nu}R(\lambda\pm i0)\langle r\rangle^{-\nu}\right|\right|_{L^2(M)\to L^2(M)}\lesssim\lambda^{-1/2}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

[Robert-Tamura,] [C. Gérard-Martinez] , [Vasy-Zworski]

High energy estimates $(\lambda \to +\infty)$ depend on the behavior of the geodesic flow ϕ^t

• Worst case: general case (everywhere below $\nu > 1/2$)

$$\left|\left|\langle r\rangle^{-\nu}R(\lambda\pm i0)\langle r\rangle^{-\nu}\right|\right|_{L^{2}(M)\to L^{2}(M)}\lesssim e^{C\lambda^{1/2}}$$

[Burq , Cardoso-Vodev]

Best case: non trapping geodesic flow

$$\left|\left|\langle r\rangle^{-\nu}R(\lambda\pm i0)\langle r\rangle^{-\nu}\right|\right|_{L^{2}(M)\to L^{2}(M)}\lesssim\lambda^{-1/2}$$

[Robert-Tamura,] [C. Gérard-Martinez] , [Vasy-Zworski]

Rem: this estimate is equivalent to the non trapping condition [Wang]

 Intermediate cases: for "weak hyperbolic trapping" (hyperbolic trapping with negative topological pressure)

$$\left|\left|\langle r
ight
angle^{-
u} R(\lambda\pm i0)\langle r
ight
angle^{-
u}
ight|
ight|_{L^2(M) o L^2(M)}\lesssim\lambda^{-1/2}\log\lambda$$

[Christianson, Datchev, Nonnenmacher-Zworski] (+ [Ikawa] for obstacles)

High energy estimates $(\lambda \to +\infty)$ depend on the behavior of the geodesic flow ϕ^t

• Worst case: general case (everywhere below $\nu > 1/2$)

$$\left|\left|\langle r\rangle^{-\nu}R(\lambda\pm i0)\langle r\rangle^{-\nu}\right|\right|_{L^{2}(M)\to L^{2}(M)}\lesssim e^{C\lambda^{1/2}}$$

[Burq , Cardoso-Vodev]

Best case: non trapping geodesic flow

$$\left|\left|\langle r\rangle^{-\nu}R(\lambda\pm i0)\langle r\rangle^{-\nu}\right|\right|_{L^{2}(M)\to L^{2}(M)}\lesssim\lambda^{-1/2}$$

[Robert-Tamura,] [C. Gérard-Martinez] , [Vasy-Zworski]

Rem: this estimate is equivalent to the non trapping condition [Wang]

 Intermediate cases: for "weak hyperbolic trapping" (hyperbolic trapping with negative topological pressure)

$$\left|\left|\langle r\rangle^{-
u}R(\lambda\pm i0)\langle r\rangle^{-
u}\right|\right|_{L^{2}(M)\to L^{2}(M)}\lesssim\lambda^{-1/2}\log\lambda$$

[Christianson, Datchev, Nonnenmacher-Zworski] (+ [Ikawa] for obstacles) For certain surfaces of revolution

$$\left|\left|\langle r\rangle^{-\nu}R(\lambda\pm i0)\langle r\rangle^{-\nu}\right|\right|_{L^2(M)\to L^2(M)}\lesssim \lambda^{\kappa}$$

[Christianson-Wunsch]

High energy estimates $(\lambda \to +\infty)$ depend on the behavior of the geodesic flow ϕ^t

• Worst case: general case (everywhere below $\nu > 1/2$)

$$\left|\left|\langle r\rangle^{-\nu}R(\lambda\pm i0)\langle r\rangle^{-\nu}\right|\right|_{L^{2}(M)\to L^{2}(M)}\lesssim e^{C\lambda^{1/2}}$$

[Burq , Cardoso-Vodev]

Best case: non trapping geodesic flow

$$\left|\left|\langle r\rangle^{-\nu}R(\lambda\pm i0)\langle r\rangle^{-\nu}\right|\right|_{L^{2}(M)\to L^{2}(M)}\lesssim\lambda^{-1/2}$$

[Robert-Tamura,] [C. Gérard-Martinez] , [Vasy-Zworski]

Rem: this estimate is equivalent to the non trapping condition [Wang]

 Intermediate cases: for "weak hyperbolic trapping" (hyperbolic trapping with negative topological pressure)

$$\left|\left|\langle r\rangle^{-
u}R(\lambda\pm i0)\langle r\rangle^{-
u}\right|\right|_{L^2(M) o L^2(M)}\lesssim\lambda^{-1/2}\log\lambda$$

[Christianson, Datchev, Nonnenmacher-Zworski] (+ [Ikawa] for obstacles) For certain surfaces of revolution

$$\left|\left|\langle r\rangle^{-\nu}R(\lambda\pm i0)\langle r\rangle^{-\nu}\right|\right|_{L^{2}(M)\to L^{2}(M)}\lesssim\lambda^{\kappa}$$

[Christianson-Wunsch]

Partial converse for trapping manifolds: if there are trapped geodesics

$$\left|\left|\langle r\rangle^{-\nu}R(\lambda\pm i0)\langle r\rangle^{-\nu}\right|\right|_{L^{2}(M)\to L^{2}(M)}\gtrsim\lambda^{-1/2}\log\lambda$$

[Bony-Burq-Ramond]

◆□ ▶ < 圖 ▶ < 圖 ▶ < 圖 ▶ < 圖 • 의 Q @</p>

Low energy estimates $(\lambda \rightarrow 0)$

Low energy estimates $(\lambda \to 0)$ In dimension $n \ge 3$, if $\nu_1, \nu_2 > 1/2$ and $\nu_1 + \nu_2 > 2$ $||\langle r \rangle^{-\nu_1} R(\lambda \pm i0) \langle r \rangle^{-\nu_2} ||_{L^2(M) \to L^2(M)} \lesssim 1$

[Bony-Hafner]



Low energy estimates $(\lambda \to 0)$ In dimension $n \ge 3$, if $\nu_1, \nu_2 > 1/2$ and $\nu_1 + \nu_2 > 2$ $||\langle r \rangle^{-\nu_1} R(\lambda \pm i0) \langle r \rangle^{-\nu_2} ||_{L^2(M) \to L^2(M)} \lesssim 1$

[Bony-Hafner]

Sharp version:

$$\left|\left|\langle r
ight
angle^{-1} R(\lambda\pm i0)\langle r
angle^{-1}
ight|\right|_{L^2(M)
ightarrow L^2(M)}\lesssim 1$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

[B.-Royer]

Low energy estimates ($\lambda \rightarrow 0$) In dimension $n \ge 3$, if $\nu_1, \nu_2 > 1/2$ and $\nu_1 + \nu_2 > 2$

$$\left|\left|\langle r\rangle^{-
u_1}R(\lambda\pm i0)\langle r\rangle^{-
u_2}\right|\right|_{L^2(M) o L^2(M)}\lesssim 1$$

[Bony-Hafner]

Sharp version:

►

$$\left|\left|\langle r\rangle^{-1}R(\lambda\pm i0)\langle r\rangle^{-1}\right|\right|_{L^{2}(M)\rightarrow L^{2}(M)}\lesssim 1$$

[B.-Royer]

Robust estimates for powers

$$\left|\left|\langle\lambda^{\frac{1}{2}}r\rangle^{-k}(\lambda^{-1}P-1\pm i0)^{-k}\langle\lambda^{\frac{1}{2}}r\rangle^{-k}\right|\right|_{L^{2}(M)\to L^{2}(M)}\lesssim 1$$

[B.-Royer]

consequence on time decay

$$\left|\left|\langle\lambda^{\frac{1}{2}}r\rangle^{-k}\varphi(\lambda^{-1}P)e^{-itP}\langle\lambda^{\frac{1}{2}}r\rangle^{-k}\right|\right|_{L^{2}(M)\to L^{2}(M)}\lesssim\langle\lambda t\rangle^{1-k}$$

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

Several results for local in time estimates

(ロ)、(型)、(E)、(E)、 E) の(の)

► For general manifolds:

Several results for local in time estimates

> For general manifolds: estimates with loss of derivatives

 $||e^{i \cdot P} u_0||_{L^p([-T,T],L^q)} \lesssim_T ||u_0||_{H^{1/p}(M)}$

Several results for local in time estimates

▶ For general manifolds: estimates with loss of derivatives

$$||e^{i \cdot P} u_0||_{L^p([-T,T],L^q)} \lesssim_T ||u_0||_{H^{1/p}(M)} := ||\langle -\Delta_G \rangle^{1/2p} u_0||_{L^2}$$

(ロ)、(型)、(E)、(E)、 E) の(の)

[Burq-Gérard-Tzvetkov]

Several results for local in time estimates

▶ For general manifolds: estimates with loss of derivatives

$$||e^{i \cdot P} u_0||_{L^p([-T,T],L^q)} \lesssim_T ||u_0||_{H^{1/p}(M)} := \left|\left|\langle -\Delta_G \rangle^{1/2p} u_0\right|\right|_{L^2}$$

(ロ)、(型)、(E)、(E)、 E) の(の)

[Burq-Gérard-Tzvetkov]

• For non trapping asymptotically flat manifolds:

Several results for local in time estimates

► For general manifolds: estimates with loss of derivatives

$$||e^{j \cdot P} u_0||_{L^p([-T,T],L^q)} \lesssim_T ||u_0||_{H^{1/p}(M)} := \left|\left|\langle -\Delta_G \rangle^{1/2p} u_0\right|\right|_{L^2}$$

[Burq-Gérard-Tzvetkov]

• For non trapping asymptotically flat manifolds:

$$||e^{i \cdot P} u_0||_{L^p([-T,T],L^q)} \lesssim_T ||u_0||_{L^2}$$

[Staffilani-Tataru], [Robbiano-Zuily], [B.-Tzvetkov], [Hassell-Tao-Wunsch]

Several results for local in time estimates

► For general manifolds: estimates with loss of derivatives

$$||e^{j \cdot P} u_0||_{L^p([-T,T],L^q)} \lesssim_T ||u_0||_{H^{1/p}(M)} := \left|\left|\langle -\Delta_G \rangle^{1/2p} u_0\right|\right|_{L^2}$$

[Burq-Gérard-Tzvetkov]

► For non trapping asymptotically flat manifolds:

$$||e^{i \cdot P} u_0||_{L^p([-T,T],L^q)} \lesssim_T ||u_0||_{L^2}$$

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

[Staffilani-Tataru], [Robbiano-Zuily], [B.-Tzvetkov], [Hassell-Tao-Wunsch]

> For asymptotically flat manifolds with small hyperbolic trapped set

Several results for local in time estimates

> For general manifolds: estimates with loss of derivatives

$$||e^{j \cdot P} u_0||_{L^p([-T,T],L^q)} \lesssim_T ||u_0||_{H^{1/p}(M)} := \left|\left|\langle -\Delta_G \rangle^{1/2p} u_0\right|\right|_{L^2}$$

[Burq-Gérard-Tzvetkov]

► For non trapping asymptotically flat manifolds:

$$||e^{i \cdot P} u_0||_{L^p([-T,T],L^q)} \lesssim T ||u_0||_{L^2}$$

[Staffilani-Tataru], [Robbiano-Zuily], [B.-Tzvetkov], [Hassell-Tao-Wunsch]

> For asymptotically flat manifolds with small hyperbolic trapped set

$$||e^{i \cdot P} u_0||_{L^p([-T,T],L^q)} \lesssim_T ||u_0||_{L^2}$$

[Burq-Guillarmou-Hassell]

Intuition (non trapping case):

▶ Inside a compact set K, combine

$$||\mathbf{1}_{K}e^{i\cdot P}u_{0}||_{L^{2}([-T,T],L^{2^{*}})} \lesssim \tau ||u_{0}||_{H^{1/2}(M)} \qquad \text{and} ||\mathbf{1}_{K}e^{i\cdot P}v_{0}||_{L^{2}([-T,T],H^{1/2})} \lesssim \tau ||v_{0}||_{L^{2}([-T,T],H^{1/2})} \leq \tau ||v_{0}||_{L^{2}([-T,T],H^{1/2})} < \tau ||$$

Outside a compact set: use that the geometry is close to a nice model (...)

Few about global in time estimates (partially due to the low energy analysis)

 Tataru , Tataru-Marzuola-Metcalfe: asymptotically euclidean case, allow relatively weak trapping at infinity

Hassell-Zhang:

Few about global in time estimates (partially due to the low energy analysis)

 Tataru , Tataru-Marzuola-Metcalfe: asymptotically euclidean case, allow relatively weak trapping at infinity

(ロ)、(型)、(E)、(E)、 E) の(の)

Hassell-Zhang: non trapping assumption,

Few about global in time estimates (partially due to the low energy analysis)

 Tataru , Tataru-Marzuola-Metcalfe: asymptotically euclidean case, allow relatively weak trapping at infinity

▶ Hassell-Zhang: non trapping assumption, special type of conical ends

Results (joint with H. Mizutani)

▲□▶ ▲圖▶ ▲圖▶ ▲圖▶ → 圖 - 釣�?

Results (joint with H. Mizutani)

Let $f_0 \in C_0^\infty(\mathbb{R})$ be such that $f_0 = 1$ near 0.

<□ > < @ > < E > < E > E のQ @

Results (joint with H. Mizutani)

Let $f_0 \in C_0^{\infty}(\mathbb{R})$ be such that $f_0 = 1$ near 0. Theorem 1 (low frequency) If $n \ge 3$ and (p, q) is admissible

 $||f_0(P)e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))} \leq C||u_0||_{L^2(M)}.$

・ロト・日本・モト・モート ヨー うへで
Let $f_0 \in C_0^{\infty}(\mathbb{R})$ be such that $f_0 = 1$ near 0. Theorem 1 (low frequency) If $n \ge 3$ and (p, q) is admissible

 $||f_0(P)e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))} \leq C||u_0||_{L^2(M)}.$

・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・

Theorem 2 (high frequency at infinity)

Let $f_0 \in C_0^{\infty}(\mathbb{R})$ be such that $f_0 = 1$ near 0. Theorem 1 (low frequency) If $n \ge 3$ and (p, q) is admissible

$$||f_0(P)e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))} \leq C||u_0||_{L^2(M)}.$$

Theorem 2 (high frequency at infinity) Assuming $n \ge 2$ and that $R(\lambda \pm i0)$ grows at most polynomially in λ ,

Let $f_0 \in C_0^{\infty}(\mathbb{R})$ be such that $f_0 = 1$ near 0. Theorem 1 (low frequency) If $n \ge 3$ and (p, q) is admissible

$$||f_0(P)e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))} \leq C||u_0||_{L^2(M)}$$

Theorem 2 (high frequency at infinity) Assuming $n \ge 2$ and that $R(\lambda \pm i0)$ grows at most polynomially in λ , there exists a compact set $K \Subset M$ such that for any (p,q) admissible

$$||\mathbf{1}_{M\setminus K}(1-f_0)(P)e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))} \leq C||u_0||_{L^2(M)}.$$

Let $f_0 \in C_0^{\infty}(\mathbb{R})$ be such that $f_0 = 1$ near 0. Theorem 1 (low frequency) If $n \ge 3$ and (p, q) is admissible

$$||f_0(P)e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))} \leq C||u_0||_{L^2(M)}$$

Theorem 2 (high frequency at infinity) Assuming $n \ge 2$ and that $R(\lambda \pm i0)$ grows at most polynomially in λ , there exists a compact set $K \Subset M$ such that for any (p,q) admissible

$$||\mathbf{1}_{M\setminus K}(1-f_0)(P)e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))} \leq C||u_0||_{L^2(M)}.$$

Theorem 3 (global space-time estimates without loss of derivatives)

Let $f_0 \in C_0^{\infty}(\mathbb{R})$ be such that $f_0 = 1$ near 0. Theorem 1 (low frequency) If $n \ge 3$ and (p, q) is admissible

$$||f_0(P)e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))} \leq C||u_0||_{L^2(M)}$$

Theorem 2 (high frequency at infinity) Assuming $n \ge 2$ and that $R(\lambda \pm i0)$ grows at most polynomially in λ , there exists a compact set $K \Subset M$ such that for any (p,q) admissible

$$||\mathbf{1}_{M\setminus K}(1-f_0)(P)e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))} \leq C||u_0||_{L^2(M)}$$

Theorem 3 (global space-time estimates without loss of derivatives) If $n \ge 3$ and the trapping is hyperbolic with negative pressure,

Let $f_0 \in C_0^{\infty}(\mathbb{R})$ be such that $f_0 = 1$ near 0. Theorem 1 (low frequency) If $n \ge 3$ and (p, q) is admissible

$$||f_0(P)e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))} \leq C||u_0||_{L^2(M)}$$

Theorem 2 (high frequency at infinity) Assuming $n \ge 2$ and that $R(\lambda \pm i0)$ grows at most polynomially in λ , there exists a compact set $K \Subset M$ such that for any (p,q) admissible

$$||\mathbf{1}_{M\setminus K}(1-f_0)(P)e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))} \leq C||u_0||_{L^2(M)}$$

Theorem 3 (global space-time estimates without loss of derivatives) If $n \ge 3$ and the trapping is hyperbolic with negative pressure, then for (p,q) admissible

$$||e^{-i \cdot P} u_0||_{L^p(\mathbb{R}; L^q(M))} \le C||u_0||_{L^2(M)}$$

Let $f_0 \in C_0^{\infty}(\mathbb{R})$ be such that $f_0 = 1$ near 0. Theorem 1 (low frequency) If $n \ge 3$ and (p, q) is admissible

$$||f_0(P)e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))} \leq C||u_0||_{L^2(M)}$$

Theorem 2 (high frequency at infinity) Assuming $n \ge 2$ and that $R(\lambda \pm i0)$ grows at most polynomially in λ , there exists a compact set $K \Subset M$ such that for any (p,q) admissible

$$||\mathbf{1}_{M\setminus K}(1-f_0)(P)e^{-i\cdot P}u_0||_{L^p(\mathbb{R};L^q(M))} \leq C||u_0||_{L^2(M)}$$

Theorem 3 (global space-time estimates without loss of derivatives) If $n \ge 3$ and the trapping is hyperbolic with negative pressure, then for (p,q) admissible

$$||e^{-i \cdot P} u_0||_{L^p(\mathbb{R}; L^q(M))} \leq C||u_0||_{L^2(M)}$$

Theorem 4 (nonlinear scattering) Under the assumptions of Theorem 3, the L^2 critical equation

$$i\partial_t u - Pu = \sigma |u|^{\frac{4}{n}} u, \qquad u_{|t=0} = u_0, \qquad \sigma = \pm 1,$$

with $||u_0||_{L^2}\ll 1,$ has a unique solution in (a subspace of) $C(\mathbb{R},L^2)\cap L^{2+\frac{4}{n}}(\mathbb{R}\times M)$ and

$$||u(t)-e^{-itP}u_{\pm}||_{L^2(M)} \rightarrow 0, \qquad t \rightarrow \pm \infty.$$

◆□ > ◆□ > ◆三 > ◆三 > ・三 ● のへで

Low frequency localization in the uncertainty region:

◆□ ▶ < 圖 ▶ < 圖 ▶ < 圖 ▶ < 圖 • 의 Q @</p>

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to 0,$ how to prove

$$\int_{\mathbb{R}} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2(\mathbb{R}; L^{2^*})}^2 dt C \left| \left| f(P/\epsilon^2) u_0 \right| \right|_{L^2}^2$$

◆□ ▶ < 圖 ▶ < 圖 ▶ < 圖 ▶ < 圖 • 의 Q @</p>

with C independent of λ (and u_0)

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to 0,$ how to prove

$$\int_{\mathbb{R}} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2(\mathbb{R}; L^{2^*})}^2 dt C \left| \left| f(P/\epsilon^2) u_0 \right| \right|_{L^2}^2$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

with C independent of λ (and u_0)

 $\left|\left|\chi(\epsilon r)f(P/\epsilon^2)e^{itP}u_0\right|\right|_{L^{2^*}}$ \lesssim

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to 0,$ how to prove

$$\int_{\mathbb{R}} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2(\mathbb{R};L^{2^*})}^2 dt C \left| \left| f(P/\epsilon^2) u_0 \right| \right|_{L^2}^2$$

・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・
・

with C independent of λ (and u_0)

 $\left|\left|\chi(\epsilon r)f(P/\epsilon^{2})e^{itP}u_{0}\right|\right|_{L^{2^{*}}} \lesssim \left|\left|\nabla_{G}\chi(\epsilon r)f(P/\epsilon^{2})e^{itP}u_{0}\right|\right|_{L^{2}}\right|$

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to 0,$ how to prove

$$\int_{\mathbb{R}} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2(\mathbb{R}; L^{2^*})}^2 dt C \left| \left| f(P/\epsilon^2) u_0 \right| \right|_{L^2}^2$$

with C independent of λ (and u_0)

 $\left|\left|\chi(\epsilon r)f(P/\epsilon^2)e^{itP}u_0\right|\right|_{L^{2^*}} \lesssim \left|\left|\nabla_G\chi(\epsilon r)f(P/\epsilon^2)e^{itP}u_0\right|\right|_{L^2} \quad \text{(homogeneous Sobolev est.)}$

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to {\rm 0,\ how}$ to prove

$$\int_{\mathbb{R}}^{\Gamma} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2(\mathbb{R}; L^{2^*})}^2 dt C \left| \left| f(P/\epsilon^2) u_0 \right| \right|_{L^2}^2$$

with C independent of λ (and u_0)

$$\begin{aligned} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^{2^*}} &\lesssim \quad \left| \left| \nabla_G \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} & \text{(homogeneous Sobolev est.)} \\ &\lesssim \quad \left| \left| \epsilon \chi'(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \end{aligned}$$

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to {\rm 0,\ how}$ to prove

$$\int_{\mathbb{R}}^{\Gamma} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2(\mathbb{R}; L^{2^*})}^2 dt C \left| \left| f(P/\epsilon^2) u_0 \right| \right|_{L^2}^2$$

with C independent of λ (and u_0)

$$\begin{aligned} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^{2^*}} &\lesssim \quad \left| \left| \nabla_G \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} & \text{(homogeneous Sobolev est.)} \\ &\lesssim \quad \left| \left| \epsilon \chi'(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \chi(\epsilon r) \nabla_G f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \end{aligned}$$

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to {\rm 0,\ how}$ to prove

$$\int_{\mathbb{R}}^{r} ||\chi(\epsilon r)f(P/\epsilon^{2})e^{itP}u_{0}||_{L^{2}(\mathbb{R};L^{2^{*}})}^{2}dtC||f(P/\epsilon^{2})u_{0}||_{L^{2}}^{2}$$

with C independent of λ (and u_0)

$$\begin{aligned} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^{2*}} &\lesssim \left| \left| \nabla_G \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} & \text{(homogeneous Sobolev est.)} \\ &\lesssim \left| \left| \epsilon \chi'(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \chi(\epsilon r) \nabla_G f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \end{aligned}$$

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to {\rm 0,\ how}$ to prove

$$\int_{\mathbb{R}}^{r} ||\chi(\epsilon r)f(P/\epsilon^{2})e^{itP}u_{0}||_{L^{2}(\mathbb{R};L^{2^{*}})}^{2}dtC||f(P/\epsilon^{2})u_{0}||_{L^{2}}^{2}$$

with C independent of λ (and u_0)

$$\begin{aligned} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^{2*}} &\lesssim \quad \left| \left| \nabla_G \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} & \text{(homogeneous Sobolev est.)} \\ &\lesssim \quad \left| \left| \epsilon \chi'(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \chi(\epsilon r) \nabla_G f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle \epsilon r \rangle^{-1} P^{\frac{1}{2}} \tilde{f}(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \end{aligned}$$

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to {\rm 0,\ how}$ to prove

$$\int_{\mathbb{R}} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2(\mathbb{R}; L^{2^*})}^2 dt C \left| \left| f(P/\epsilon^2) u_0 \right| \right|_{L^2}^2 dt C \right| dt C dt$$

with C independent of λ (and u_0)

$$\begin{aligned} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^{2^*}} &\lesssim \quad \left| \left| \nabla_G \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} & \text{(homogeneous Sobolev est.)} \\ &\lesssim \quad \left| \left| \epsilon \chi'(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \chi(\epsilon r) \nabla_G f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle \epsilon r \rangle^{-1} P^{\frac{1}{2}} \tilde{f}(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \end{aligned}$$

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to {\rm 0,\ how}$ to prove

$$\int_{\mathbb{R}} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2(\mathbb{R}; L^{2^*})}^2 dt C \left| \left| f(P/\epsilon^2) u_0 \right| \right|_{L^2}^2 dt C \right| dt C dt$$

with C independent of λ (and u_0)

$$\begin{aligned} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^{2^*}} &\lesssim \quad \left| \left| \nabla_G \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} & \text{(homogeneous Sobolev est.)} \\ &\lesssim \quad \left| \left| \epsilon \chi'(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \chi(\epsilon r) \nabla_G f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle \epsilon r \rangle^{-1} \frac{P^{\frac{1}{2}}}{\tilde{f}} (P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle r \rangle^{-1} \frac{\tilde{f}}{\tilde{f}} (P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \end{aligned}$$

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to 0,$ how to prove

$$\int_{\mathbb{R}} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2(\mathbb{R}; L^{2^*})}^2 dt C \left| \left| f(P/\epsilon^2) u_0 \right| \right|_{L^2}^2 dt C \right| dt C dt$$

with C independent of λ (and u_0)

$$\begin{aligned} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^{2*}} &\lesssim \left| \left| \nabla_G \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} & \text{(homogeneous Sobolev est.)} \\ &\lesssim \left| \left| \epsilon \chi'(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \chi(\epsilon r) \nabla_G f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle \epsilon r \rangle^{-1} \frac{P^{\frac{1}{2}}}{\tilde{f}} (P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle r \rangle^{-1} \tilde{f}(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \end{aligned}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへで

où $\widetilde{f}, \widetilde{\widetilde{f}} \in C_0^\infty(0, +\infty).$

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to 0,$ how to prove

$$\int_{\mathbb{R}} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2(\mathbb{R}; L^{2^*})}^2 dt C \left| \left| f(P/\epsilon^2) u_0 \right| \right|_{L^2}^2 dt C \right| dt C dt$$

with C independent of λ (and u_0)

$$\begin{aligned} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^{2^*}} &\lesssim \quad \left| \left| \nabla_G \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} & \text{(homogeneous Sobolev est.)} \\ &\lesssim \quad \left| \left| \epsilon \chi'(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \chi(\epsilon r) \nabla_G f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle \epsilon r \rangle^{-1} \frac{P^{\frac{1}{2}}}{\tilde{f}} (P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle r \rangle^{-1} \tilde{f} (P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \end{aligned}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

où $\tilde{f}, \tilde{\tilde{f}} \in C_0^\infty(0, +\infty)$. One concludes by mean of an optimally weighted resolvent inequality [B-Royer, 2015]

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to 0,$ how to prove

$$\int_{\mathbb{R}}^{r} ||\chi(\epsilon r)f(P/\epsilon^{2})e^{itP}u_{0}||_{L^{2}(\mathbb{R};L^{2^{*}})}^{2}dtC||f(P/\epsilon^{2})u_{0}||_{L^{2}}^{2}$$

with C independent of λ (and u_0)

$$\begin{aligned} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^{2^*}} &\lesssim \quad \left| \left| \nabla_G \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} & \text{(homogeneous Sobolev est.)} \\ &\lesssim \quad \left| \left| \epsilon \chi'(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \chi(\epsilon r) \nabla_G f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle \epsilon r \rangle^{-1} \frac{P^{\frac{1}{2}}}{\tilde{f}} (P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle r \rangle^{-1} \tilde{f} (P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \end{aligned}$$

(日) (日) (日) (日) (日) (日) (日) (日)

où $\tilde{f}, \tilde{\tilde{f}} \in C_0^\infty(0, +\infty)$. One concludes by mean of an optimally weighted resolvent inequality [B-Royer, 2015]

 $\left|\left|\langle r\rangle^{-1}f(P/\epsilon^2)e^{i\cdot P}u_0\right|\right|_{L^2(\mathbb{R};L^2)}\lesssim$

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to 0,$ how to prove

$$\int_{\mathbb{R}} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2(\mathbb{R};L^{2^*})}^2 dt C \left| \left| f(P/\epsilon^2) u_0 \right| \right|_{L^2}^2$$

with C independent of λ (and u_0)

$$\begin{aligned} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^{2^*}} &\lesssim \quad \left| \left| \nabla_G \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} & \text{(homogeneous Sobolev est.)} \\ &\lesssim \quad \left| \left| \epsilon \chi'(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \chi(\epsilon r) \nabla_G f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle \epsilon r \rangle^{-1} \frac{P^{\frac{1}{2}}}{\tilde{f}} (P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle r \rangle^{-1} \tilde{f} (P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \end{aligned}$$

où $\tilde{f}, \tilde{\tilde{f}} \in C_0^\infty(0, +\infty)$. One concludes by mean of an optimally weighted resolvent inequality [B-Royer, 2015]

$$\left|\left|\langle r\rangle^{-1}f(P/\epsilon^{2})e^{i\cdot P}u_{0}\right|\right|_{L^{2}(\mathbb{R};L^{2})} \lesssim \left(1+\sup_{|\lambda|\leq 2}\left|\left|\langle r\rangle^{-1}(P-\lambda\pm i0)^{-1}\langle r\rangle^{-1}\right|\right|_{L^{2}\to L^{2}}\right)\left|\left|u_{0}\right|\right|_{L^{2}}\right)$$

(日) (日) (日) (日) (日) (日) (日) (日)

Rem. For the localization, $(1 - \chi(\epsilon r))f(P/\epsilon^2)$,

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to 0,$ how to prove

$$\int_{\mathbb{R}} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2(\mathbb{R};L^{2^*})}^2 dt C \left| \left| f(P/\epsilon^2) u_0 \right| \right|_{L^2}^2$$

with C independent of λ (and u_0)

$$\begin{aligned} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^{2^*}} &\lesssim \quad \left| \left| \nabla_G \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} & \text{(homogeneous Sobolev est.)} \\ &\lesssim \quad \left| \left| \epsilon \chi'(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \chi(\epsilon r) \nabla_G f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle \epsilon r \rangle^{-1} \frac{P^{\frac{1}{2}}}{\tilde{f}} (P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle r \rangle^{-1} \tilde{f} (P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \end{aligned}$$

où $\tilde{f}, \tilde{\tilde{f}} \in C_0^\infty(0, +\infty)$. One concludes by mean of an optimally weighted resolvent inequality [B-Royer, 2015]

$$\left|\left|\langle r\rangle^{-1}f(P/\epsilon^2)e^{i\cdot P}u_0\right|\right|_{L^2(\mathbb{R};L^2)} \lesssim \left(1+\sup_{|\lambda|\leq 2}\left|\left|\langle r\rangle^{-1}(P-\lambda\pm i0)^{-1}\langle r\rangle^{-1}\right|\right|_{L^2\to L^2}\right)\left|\left|u_0\right|\right|_{L^2}\right)$$

Rem. For the localization, $(1 - \chi(\epsilon r))f(P/\epsilon^2)$, one has " $|\xi| \sim \epsilon$ " and " $|x| \gtrsim \epsilon^{-1}$ "

▲□▶ ▲圖▶ ▲≣▶ ▲≣▶ = 悪 = のへで

Low frequency localization in the uncertainty region:in the regime $\lambda=\epsilon^2\to 0,$ how to prove

$$\int_{\mathbb{R}} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2(\mathbb{R};L^{2^*})}^2 dt C \left| \left| f(P/\epsilon^2) u_0 \right| \right|_{L^2}^2$$

with C independent of λ (and u_0)

$$\begin{aligned} \left| \left| \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^{2^*}} &\lesssim \quad \left| \left| \nabla_G \chi(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} & \text{(homogeneous Sobolev est.)} \\ &\lesssim \quad \left| \left| \epsilon \chi'(\epsilon r) f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \chi(\epsilon r) \nabla_G f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle \epsilon r \rangle^{-1} \frac{P^{\frac{1}{2}}}{\tilde{f}} (P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \\ &\lesssim \quad \left| \left| \langle r \rangle^{-1} f(P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} + \left| \left| \langle r \rangle^{-1} \tilde{f} (P/\epsilon^2) e^{itP} u_0 \right| \right|_{L^2} \end{aligned}$$

où $\tilde{f}, \tilde{\tilde{f}} \in C_0^\infty(0, +\infty)$. One concludes by mean of an optimally weighted resolvent inequality [B-Royer, 2015]

$$\left|\left|\langle r\rangle^{-1}f(P/\epsilon^2)e^{i\cdot P}u_0\right|\right|_{L^2(\mathbb{R};L^2)} \lesssim \left(1+\sup_{|\lambda|\leq 2}\left|\left|\langle r\rangle^{-1}(P-\lambda\pm i0)^{-1}\langle r\rangle^{-1}\right|\right|_{L^2\to L^2}\right)\left|\left|u_0\right|\right|_{L^2}\right)$$

Rem. For the localization, $(1 - \chi(\epsilon r))f(P/\epsilon^2)$, one has " $|\xi| \sim \epsilon$ " and " $|x| \gtrsim \epsilon^{-1}$ " \Rightarrow no problem of uncertainty principle to use microlocal techniques

Rest of the proof

At infinity: split $f(P/\lambda)e^{itP}$ into sums of

$$T_{\lambda}(t) = L_{\lambda} f(P/\lambda) e^{itP}$$

with suitable localization operators L_{λ} , and show

 $\||\mathcal{T}_{\lambda}(t)||_{L^2 o L^2} \lesssim 1, \qquad \|\mathcal{T}_{\lambda}(t)\mathcal{T}_{\lambda}(s)||_{L^1 o L^\infty} \lesssim |t-s|^{-rac{n}{2}}$

by writing

$$\mathcal{T}_\lambda(t)\mathcal{T}_\lambda(s) = \mathsf{approximation} + \mathsf{remainder}$$

- ▶ the "approximation" is explicit enough operator to bound sharply its integral kernel by $|t s|^{-\frac{\pi}{2}}$ (dispersion bound)
- ▶ the remainder is a remainder term in a Duhamel formula in which we combine L^2 time decay/propagation estimates (for the time decay) and Sobolev estimates (to replace $L^2 \rightarrow L^2$ by $L^1 \rightarrow L^\infty$) to derive dispersion bounds.